

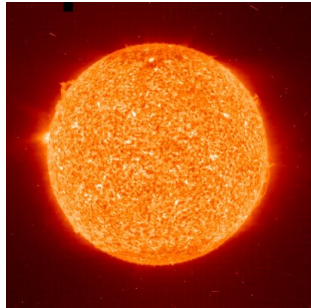
Neutrinos: Status and Perspectives

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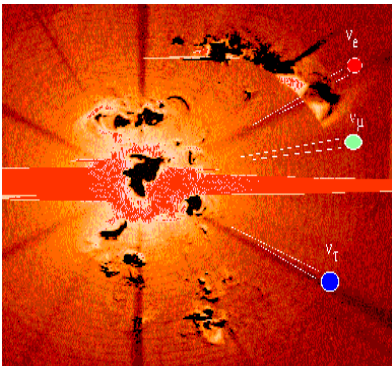


New Physics & Neutrino Sources



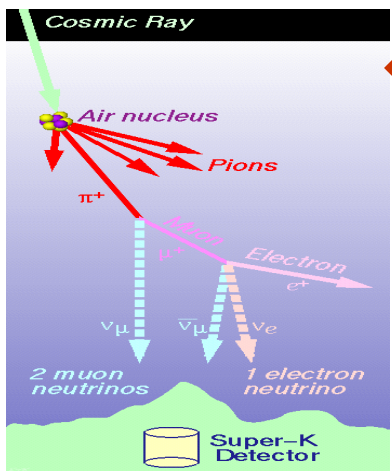
← Sun

Astronomy: →
Supernovae
GRBs
UHE ν 's



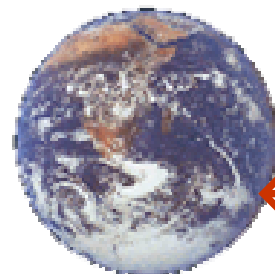
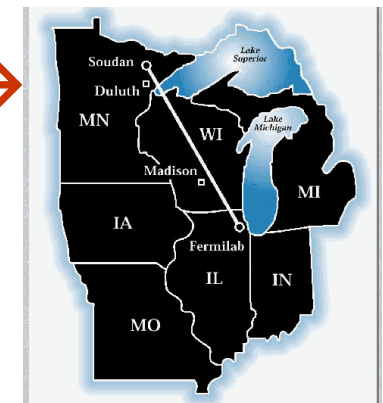
← **Cosmology**

Reactors →



← **Atmosphere**

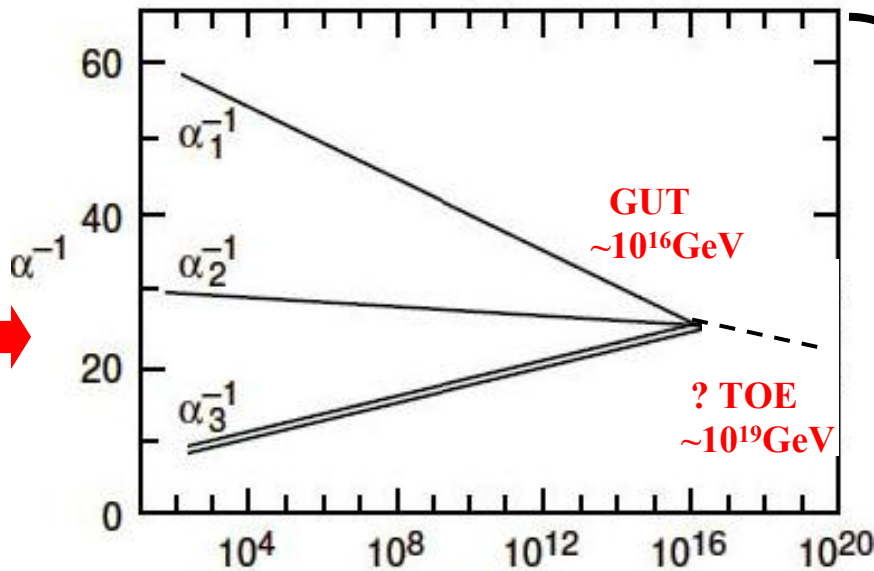
Accelerators →



← **Earth**

Different Routes Beyond the SM

gauge bosons →



experimental facts:

Dark Matter

Dark Energy

neutrino masses & mixings

baryon asymmetry $\leftrightarrow m_\nu > 0$

→ indirect tests of VHE physics

Higgs →

gauge hierarchy problem:

$$\delta m_H^2 \sim \Lambda^2$$

quarks leptons →

flavour problem: 3 generations
many parameters (m_i , mixings)
unification into GUTs

$$m_\nu = (m^D)^T M_R^{-1} m_D$$

SUSY
~TeV

$\sim \Lambda_{\text{GUT}}$
+seesaw

astrophysics
& cosmology

Extending the Standard Model

→ success of renormalizable gauge field theories in $d=4$

QED	→	QCD	→	SM
$U(1)_{em}$		$SU(3)_C$		$SU(3)_C \times SU(2)_L \times U(1)_Y$

→ symmetry, renormalizability, no anomalies

→ particle content (symmetry representations):

gauge sector – fixed by gauge group

scalar sector – must break EW symmetry, $SB \sim 2_L$

fermions – anomaly free combinations

→ different levels of SM extension...

- add further representations
- extend the gauge symmetry
- add supersymmetry
- extend/modify basic concepts: quantum fields and/or space-time

Adding Neutrino Mass Terms

1) Postulate right handed neutrino fields \rightarrow SM+

Field	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
$L_Q = \begin{pmatrix} l_u \\ l_d \end{pmatrix}$	3	2	1/3
r_u	3	1	4/3
r_d	3	1	-2/3
$L_L = \begin{pmatrix} l_\nu \\ l_e \end{pmatrix}$	1	2	-1
$r_\nu ???$	1	1	0
r_e	1	1	-2

not part of SM !

makes table more symmetric
3 right handed neutrinos?

NEW: \rightarrow 9 parameters

\rightarrow explicit fermion mass term

\rightarrow L number violation

$$\begin{array}{c} \nu_L \quad \bar{\nu}_L \quad \nu_R \\ \hline \phantom{\nu_L \quad \bar{\nu}_L \quad \nu_R} \\ \phantom{\nu_L \quad \bar{\nu}_L \quad \nu_R} \times \\ \phantom{\nu_L \quad \bar{\nu}_L \quad \nu_R} \downarrow \\ \langle \phi \rangle = v \end{array}$$

$$\begin{array}{c} \nu_R \quad \nu_R \\ \hline \times \\ \text{Majorana} \\ \cancel{\neq} \end{array}$$



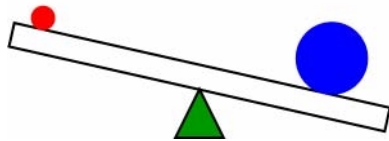
$$\begin{pmatrix} \bar{\nu}_L & \bar{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

Natural value of mass operators: scale of symmetry

$m_D \sim$ electro-weak scale

$M_R \sim$ L violation scale \leftrightarrow embedding into GUTs

See-saw mechanism (type I)



$$m_\nu = m_D M_R^{-1} m_D^T$$

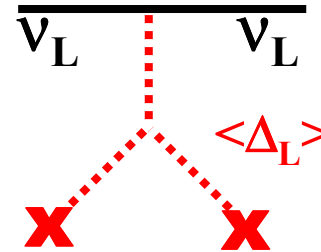
$$m_h = M_R$$

Suggestive hints:

For $m_3 \sim (\Delta m_{\text{atm}}^2)^{1/2}$, $m_D \sim$ leptons $\rightarrow M_R \sim 10^{11} - 10^{16} \text{ GeV}$
 $\rightarrow \nu$'s are **Majorana particles**, m_ν probes \sim **GUT scale physics!**
 \rightarrow smallness of m_ν \leftrightarrow high scale of L , symmetries of m_D, M_R

Other Neutrino Mass Operators

2) new Higgs triplets Δ_L :



→ left-handed Majorana mass term:

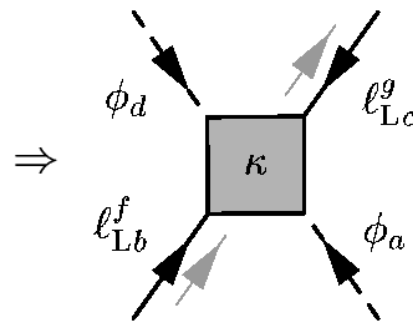
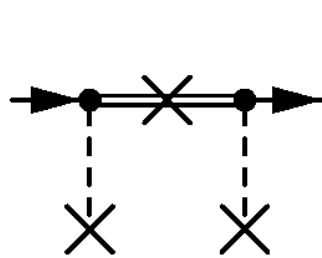
$$\rightarrow M_L \bar{L} L^c$$

3) Both ν_R and new Higgs triplets Δ_L :

→ see-saw type II

$$m_\nu = M_L - m_D M_R^{-1} m_D^T$$

4) Higher dimensional operators: $d=5, \dots$



\Leftrightarrow

$$\mathcal{L}_{mass} = \kappa \cdot \bar{\nu}_L^C \nu_L \Phi^T \Phi$$

$$\rightarrow M_L \bar{L} L^c$$

Other effective Operators Beyond the SM

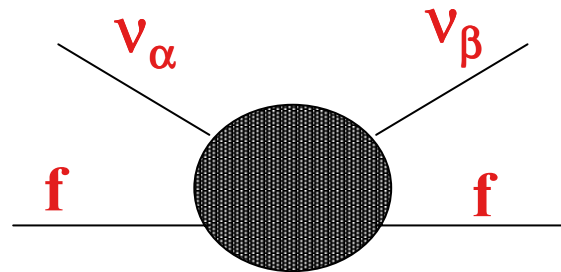
→ effects beyond 3 flavours

→ **Non Standard Interactions = NSIs** → effective 4f operators

$$\mathcal{L}_{NSI} \simeq \epsilon_{\alpha\beta} 2\sqrt{2}G_F (\bar{\nu}_{L\beta} \gamma^\rho \nu_{L\alpha}) (\bar{f}_L \gamma_\rho f_L)$$

• **integrating out heavy physics (c.f. $G_F \leftrightarrow M_W$)**

$$|\epsilon| \simeq \frac{M_W^2}{M_{NSI}^2}$$



Grossman, Bergmann+Grossman, Ota+Sato, Honda et al., Friedland+Lunardini,
Blennow+Ohlsson+Skrotzki, Huber+Valle, Huber+Schwetz+Valle,
Campanelli+Romanino, Bueno et al., Barranco+Miranda+Rashba, Kopp+ML+Ota, ...

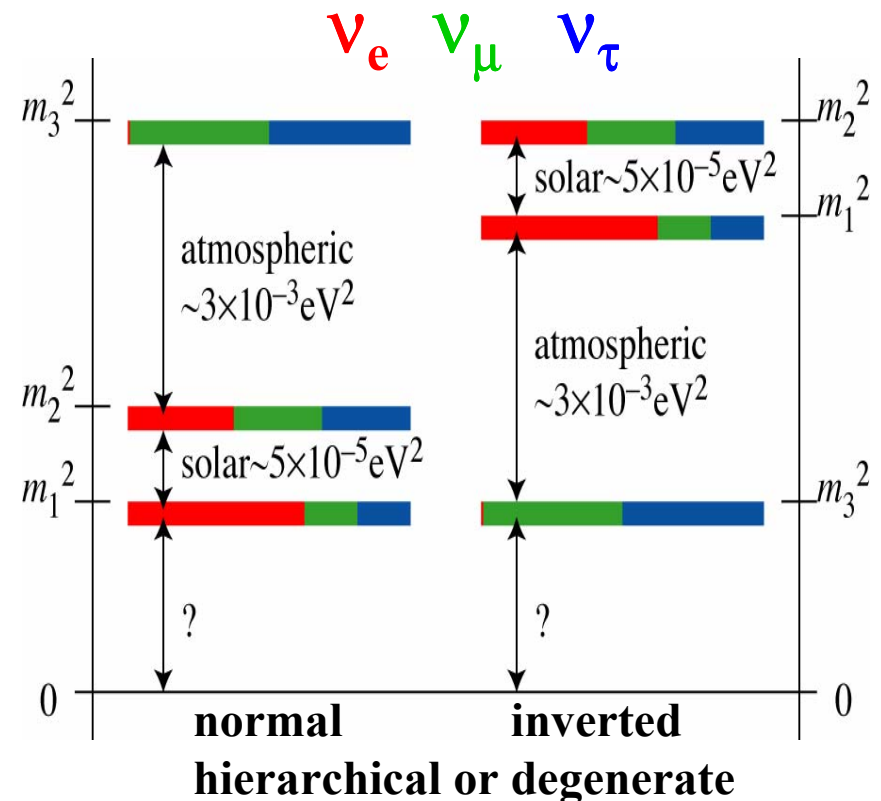
Parameters for 3 Light Neutrinos

mass & mixing parameters: m_1 , Δm^2_{21} , $|\Delta m^2_{31}|$, $\text{sign}(\Delta m^2_{31})$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \text{diag}(e^{i\alpha}, e^{i\beta}, 1)$$

questions:

- Dirac / Majorana
- mass scale: m_1
- mass ordering: $\text{sgn}(\Delta m^2_{31})$
- how small is θ_{13} , θ_{23} maximal?
- leptonic CP violation
- 3 flavour unitarity?
- why 3 generations, $d=4$, gauge group, ...



Four Methods of Mass Determination

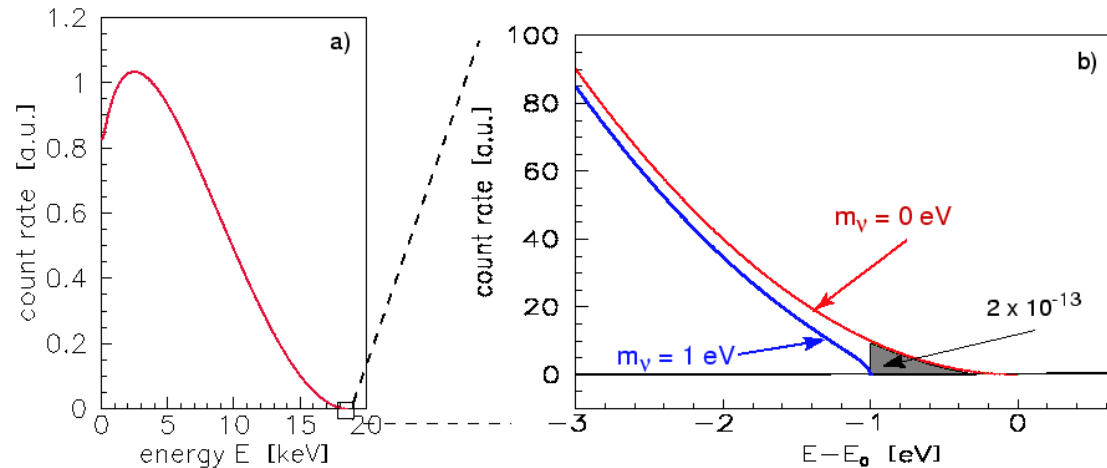
- **kinematical**
- **cosmology**
- **lepton number violation**
 \leftrightarrow **ν -less double β -decay**
- **oscillations**

Kinematical Mass Determination

Relativistic kinematics:

$$E^2 = p^2 + m^2; \quad \sum p_i^\mu = \sum p_f^\mu$$

Endpoint of decays:



Bounds:

“Elektron-Neutrino”: $m < 2.2 \text{ eV}$ (Mainz, Troitsk)

“Muon-Neutrino”: $m < 170 \text{ keV}$

“Tau-Neutrino”: $m < 15.5 \text{ MeV}$

Sensitivity \Leftrightarrow degenerate ν -spectrum

$$\Rightarrow \text{Oscillations: } \Delta m_{ij}^2 \ll m_i^2 \Rightarrow \sum m_i^2 |U_{ei}|^2 < (2.2 \text{ eV})^2$$

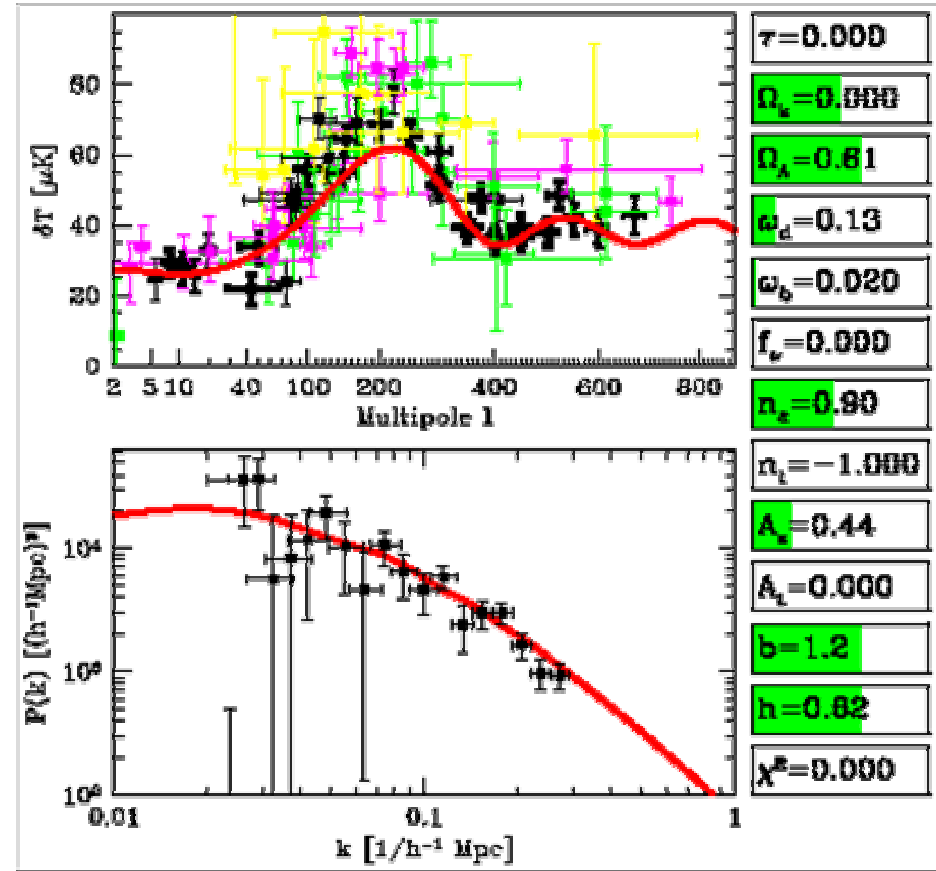
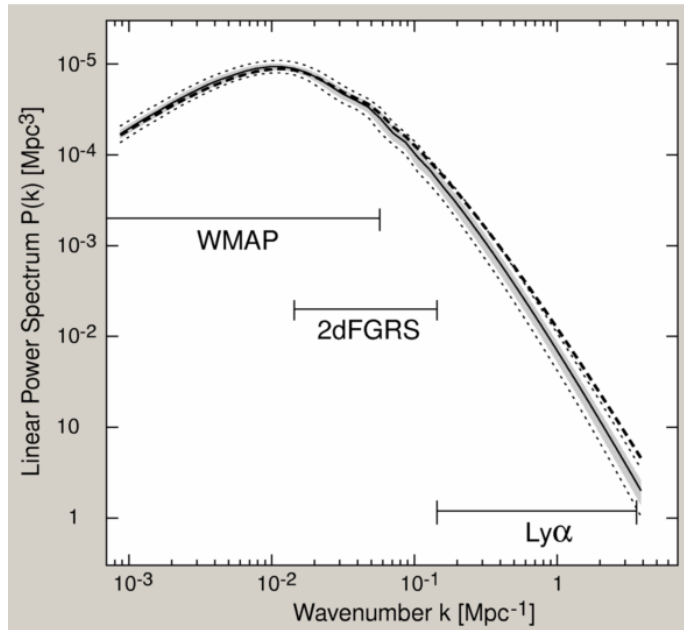
Future: KATRIN \rightarrow 0.25 eV

\leftrightarrow c.f. cosmological bounds

Cosmology and Neutrino Mass

- ν 's hot dark matter \rightarrow smears structure @ small scales

Tegmark



- WMAP+2dFGRS + Ly α +...

\rightarrow bound: $\Sigma m_\nu < 0.17 - 1.2$ eV

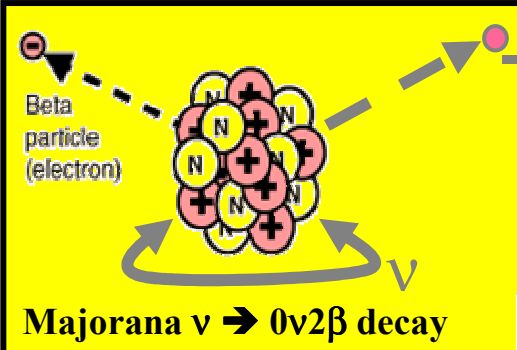
\leftrightarrow levels of systematic errors

- 3 degenerate neutrinos, conservative approach

$\rightarrow m_\nu < 0.25$ eV future improvements: \sim factor 5-10 ?

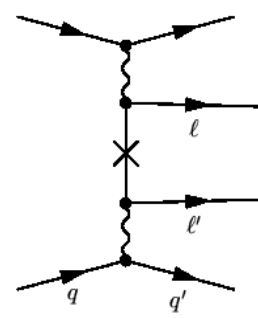
$$f_\nu = \Omega_\nu / \Omega_{\text{matter}}$$

Neutrino-less Double β -Decay



Beta particle (electron)

Majorana $\nu \rightarrow 0\nu 2\beta$ decay



$\propto |\langle m_{ee} \rangle| = |\sum m_i U_{ei}^2| \leq 0.35 \text{ eV} ?$

Heidelberg-Moscow experiment

$$m_{ee} = |m_{ee}^{(1)}| + |m_{ee}^{(2)}| \cdot e^{i\Phi_2} + |m_{ee}^{(3)}| \cdot e^{i\Phi_3}$$

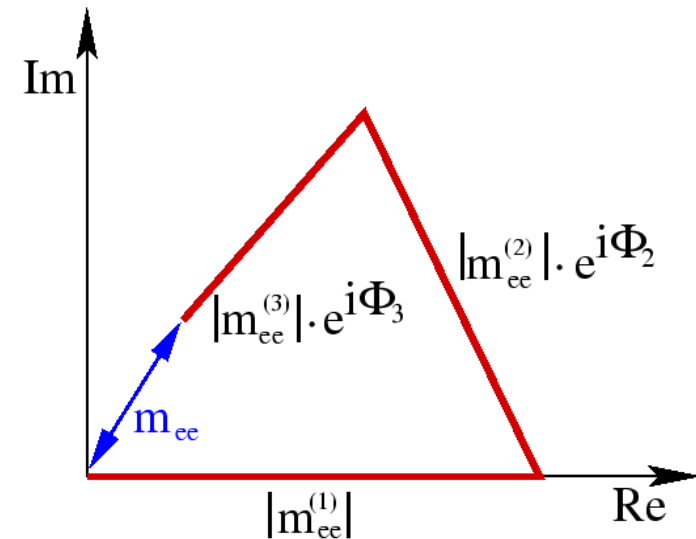
$$|m_{ee}^{(1)}| = |U_{e1}|^2 m_1$$

$$|m_{ee}^{(2)}| = |U_{e2}|^2 \sqrt{m_1^2 + \Delta m_{21}^2}$$

$$|m_{ee}^{(3)}| = |U_{e3}|^2 \sqrt{m_1^2 + \Delta m_{31}^2}$$

solar $\Rightarrow |U_{e1}|^2, |U_{e2}|^2, \Delta m_{21}^2$ atmosph. $\Rightarrow |\Delta m_{31}^2|$ CHOOZ $\Rightarrow |U_{e3}|^2 < 0.05$

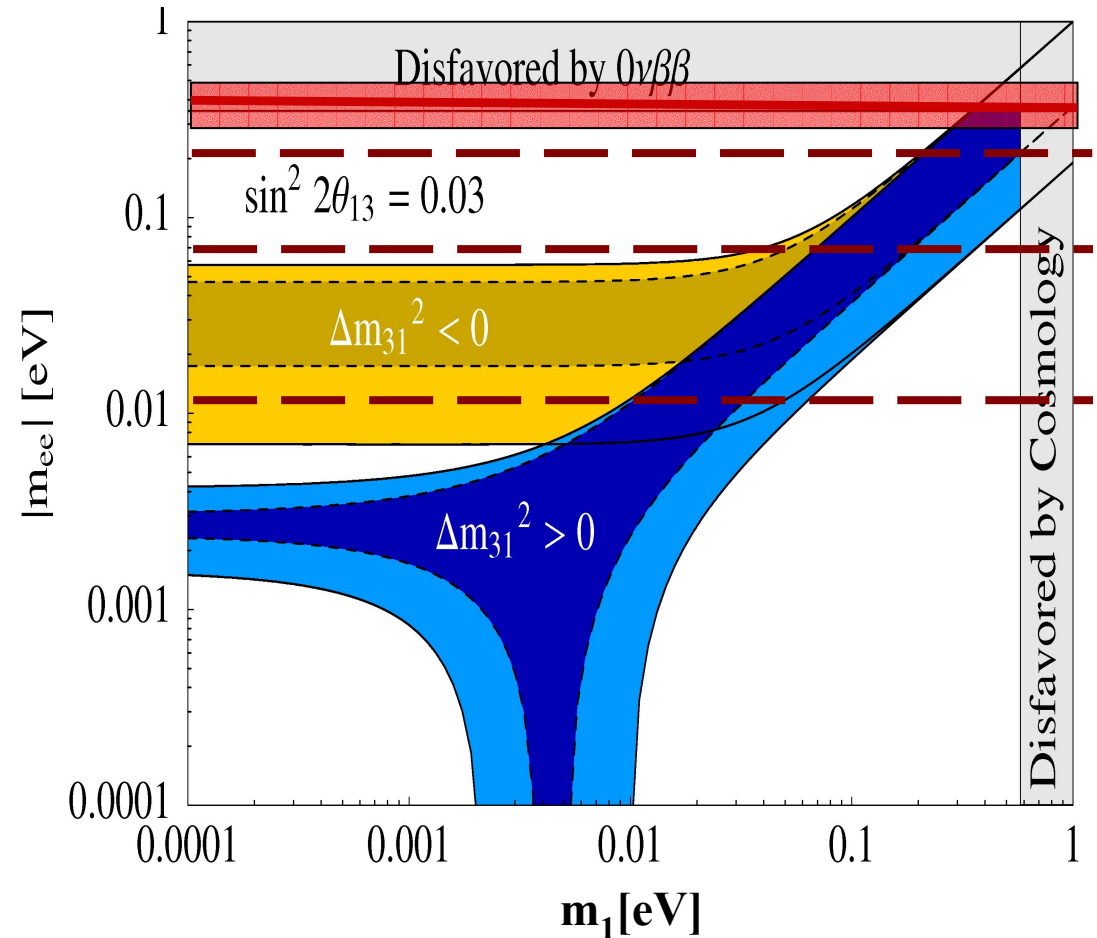
\rightarrow free parameters: $m_1, \text{sign}(\Delta m_{31}^2), \text{CP-phases } \Phi_2, \Phi_3$



Claim of part of the original Heidelberg-Moscow experiment
 \leftrightarrow cosmology \rightarrow ,tension‘

aims of new experiments:

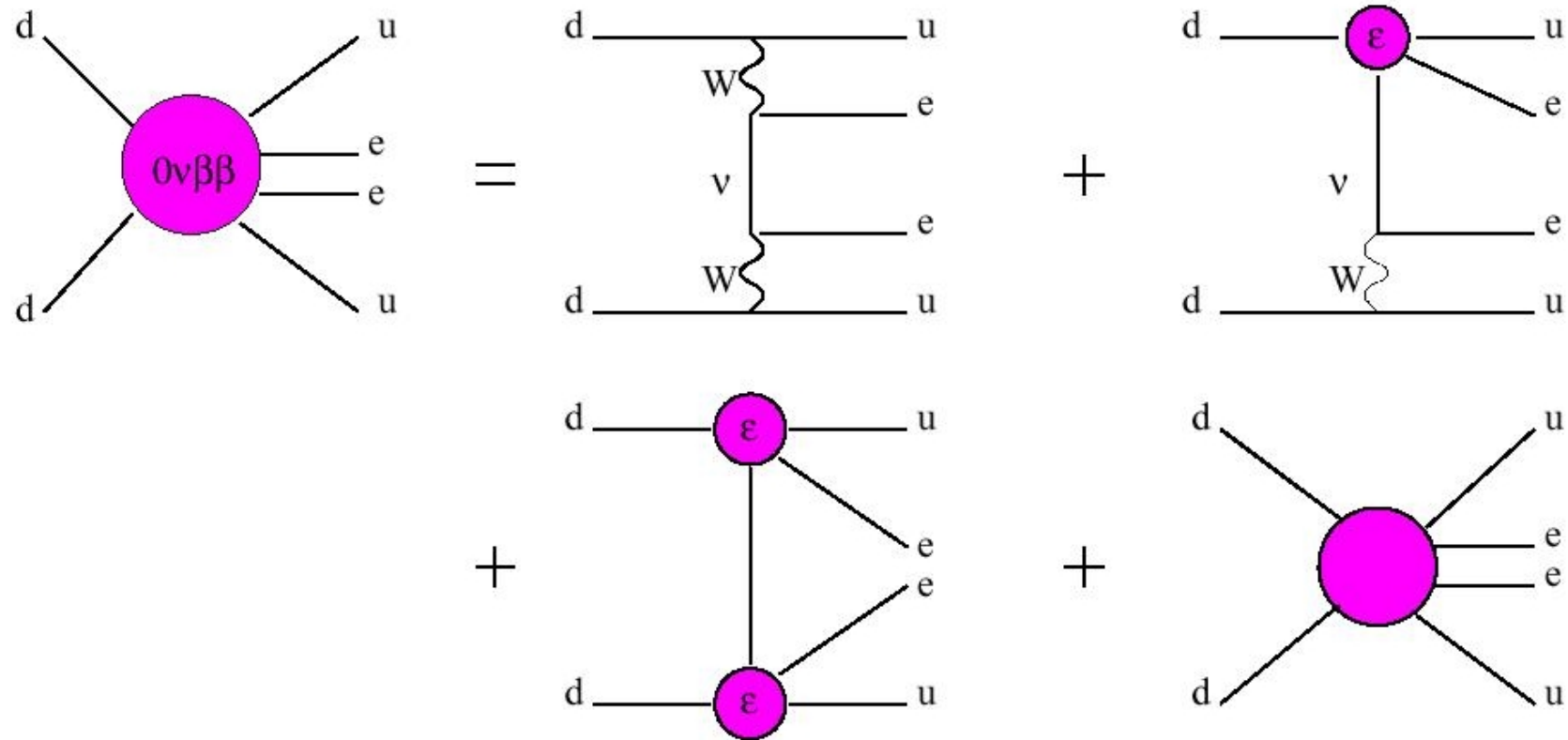
- test HM claim
- $(\Delta m_{31}^2)^{1/2} \simeq 0.05\text{eV} \pm \text{errors}$
 - \rightarrow reach **0.01eV**
 - \rightarrow CUORE
 - \rightarrow GERDA phases I, II, (III)



Limitations:

- cosmology: systematical errors \rightarrow ~factor 10?
- errors of $0\nu 2\beta$ nuclear matrix elements! ~factor 2 **theoretical** uncertainty in m_{ee}
- $\Delta m^2 > 0$ allows complete cancellation \rightarrow $0\nu 2\beta$ signal not guaranteed
- $0\nu 2\beta$ signal \rightarrow *some* lepton number violating operator

alternatives: LR, RPV-SUSY, ... \rightarrow other \cancel{L} operators \leftrightarrow NSI's



Schechter+Valle:

L violating operator \rightarrow radiative mass generation \rightarrow Majorana nature of ν 's

However: This may only be a tiny correction to a much larger Dirac mass term

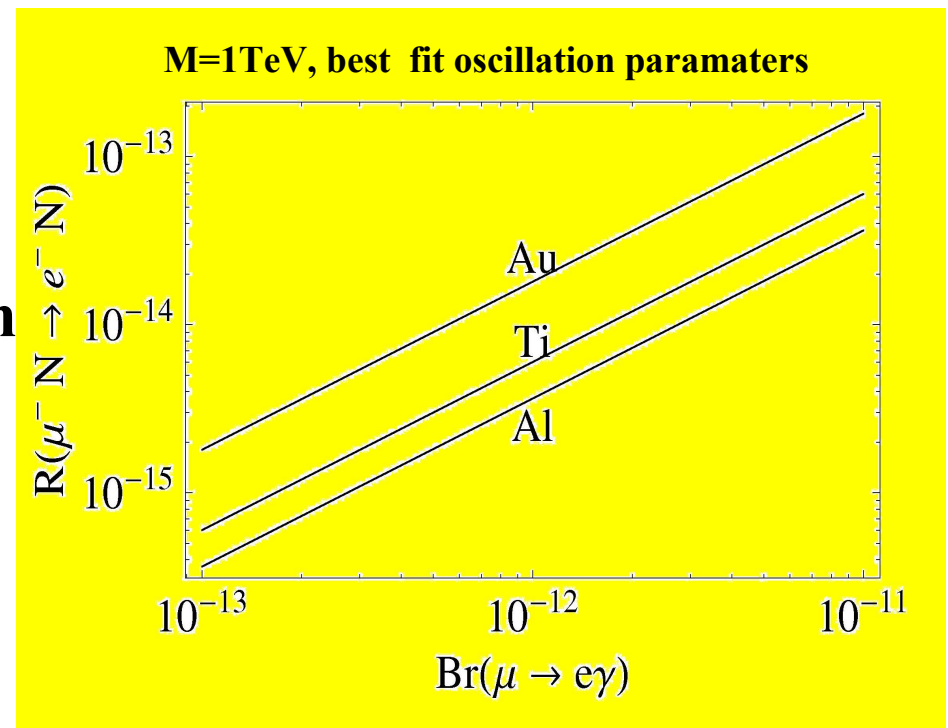
Lepton Flavour Violation

- Majorana neutrino mass terms
- ...
- **R-parity violating supersymmetry**
Hall+Kosteleck+Rabi, Borzumati+Masiero, Hisano+Tobe, Casas+Ibarra, Antusch+Arganda+Herrero+Teixeira, Joaquim+Rossi, ...

→ **LFV and leptonic CP violation can even exist for $m_\nu \rightarrow 0$**

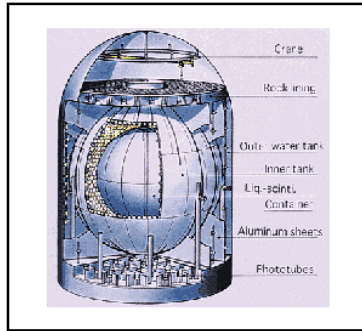
→ **e.g. modifications of correlation between $\mu^- \rightarrow e^- \gamma$ decay and nuclear $\mu^- \rightarrow e^-$ conversion**
MEG: 10^{-13}
PRISM: 10^{-18}

→ **interplay: ν 's – LFV – LHC**



Deppisch+Kosmas+Valle

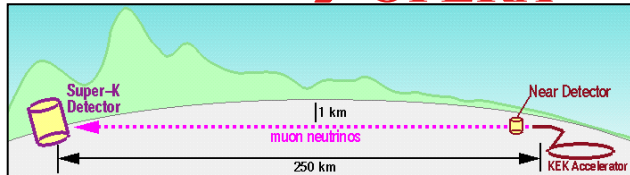
Status of Neutrino Oscillations



Reactors: KAMLAND

improved result

**Beams: K2K → MINOS
→ OPERA**



$$\Delta m_{21}^2 = (7.9 \pm 0.3) * 10^{-5} \text{ eV}^2$$

$$\tan^2 \theta_{12} = 0.39 \pm 0.05$$

$$\Delta m_{31}^2 = (2.4 \pm 0.3) * 10^{-3} \text{ eV}^2$$

$$\tan^2 \theta_{23} = 1.0 \pm 0.3$$

$$\sin^2 2\theta_{13} < 0.16 \text{ Chooz}$$

solar: GALLEX/GNO → SK, SNO

Primary neutrino source: $\rho + \rho \rightarrow \text{D} + e^+ + \nu_e$

Other sources of neutrinos:
 $e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$
 ${}^8\text{B} \rightarrow 2 {}^4\text{He} + e^+ + \nu_e$

atmospheric: Superkamiokande

Atmospheric neutrino source

LSND? → MiniBooNE

Proton beam → Water target → Pions → Muons and electrons → Neutrinos $\nu_e, \nu_\mu, \bar{\nu}_\mu$ → $\bar{\nu}_e$ detector

30 meters

- LSND not confirmed!
 → 3+2 scenarios?
 → new anomaly
 - upturn at low E?

Future Precision Oscillation Physics

Precise measurements \rightarrow 3f oscillation formulae

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \text{Majorana-CP-phases}$$

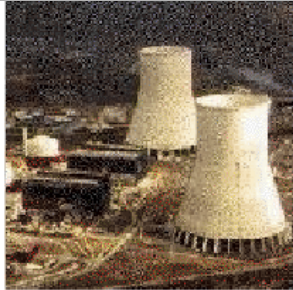
θ_{23} $S_{13} \rightarrow$ 3 flavour effects \rightarrow CP phase δ θ_{12}

Aims: \rightarrow **improved precision** of the leading 2x2 oscillations
 \rightarrow **detection of generic 3-neutrino effects: θ_{13} , CP violation**

Complication: **Matter effects \rightarrow effective parameters in matter**
 \rightarrow **expansion in small quantities θ_{13} and $a = \Delta m^2_{\text{sol}} / \Delta m^2_{\text{atm}}$**

Burguet-Castell et al. , Akhmedov et al. ...

Future Precision with Reactor Experiments



$\bar{\nu}_e$

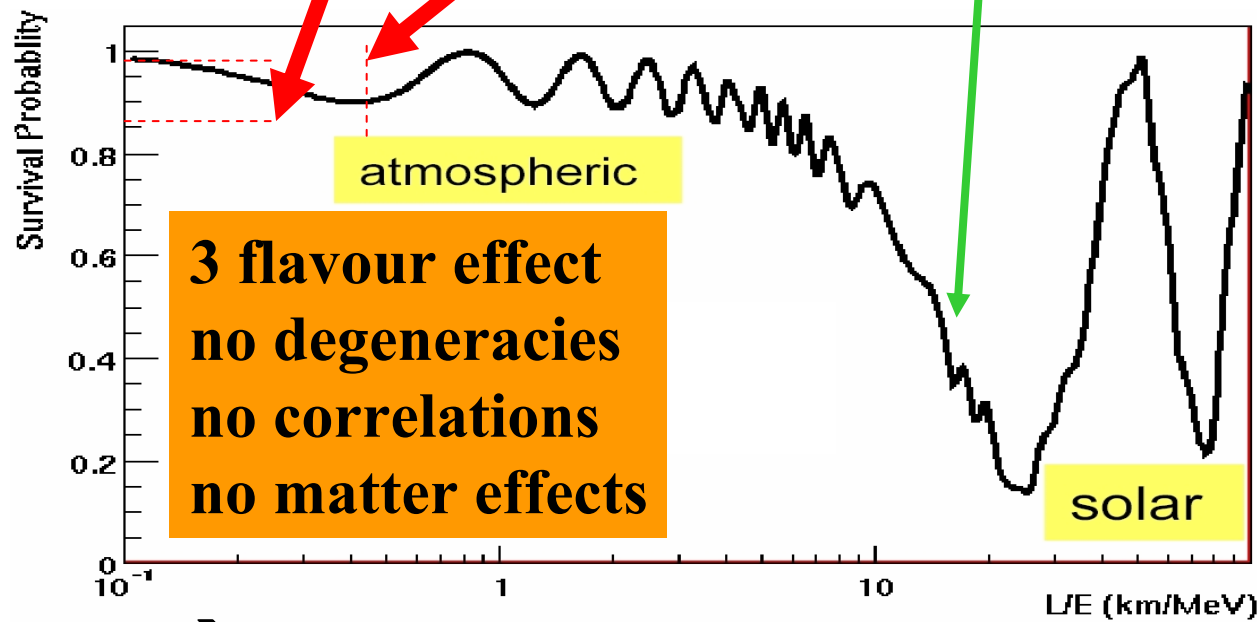
near detector (170m)

$\bar{\nu}_e$

far detector (1700m)

identical detectors → many errors cancel

$$P_{ee} \approx 1 - \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{31}^2 L}{4E_\nu} - \left(\frac{\Delta m_{21}^2 L}{4E_\nu} \right)^2 \cos^4 \theta_{13} \sin^2 2\theta_{12}$$



E=4MeV → 2km 4km 40km 80km

- Double Chooz
- Daya Bay
- Reno
- Angra

clean & precise θ_{13} measurements

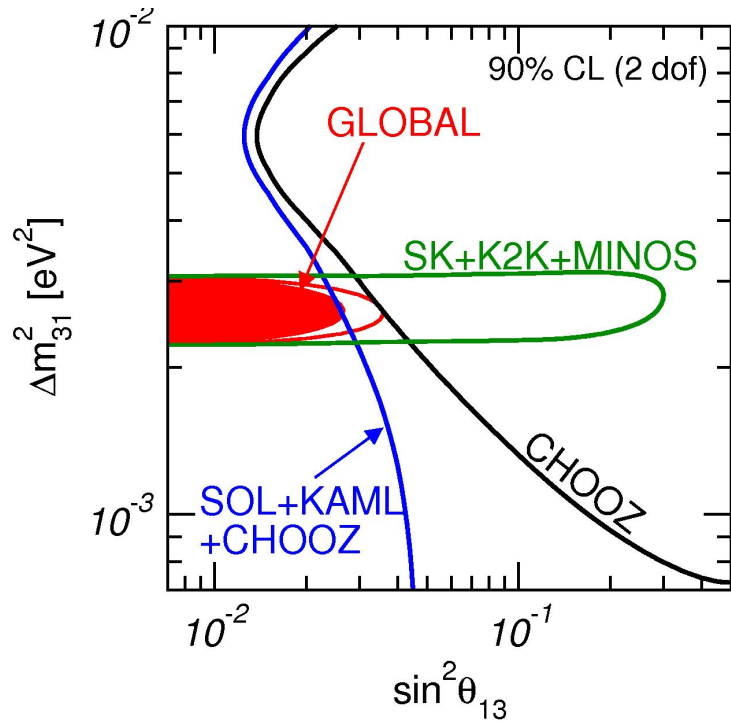
Future Precision with New Neutrino Beams

- conventional beams, superbeams
→ MINOS, CNGS, T2K, NOvA, T2H,...
- β -beams
→ pure ν_e and $\bar{\nu}_e$ beams from radioactive decays; $\gamma \simeq 100$
- neutrino factories
→ clean neutrino beams from decay of stored μ 's

$$\begin{aligned} P(\nu_e \rightarrow \nu_\mu) &\approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2((1-\hat{A})\Delta)}{(1-\hat{A})^2} \\ &\pm \sin \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\ &+ \cos \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\ &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2} \end{aligned}$$

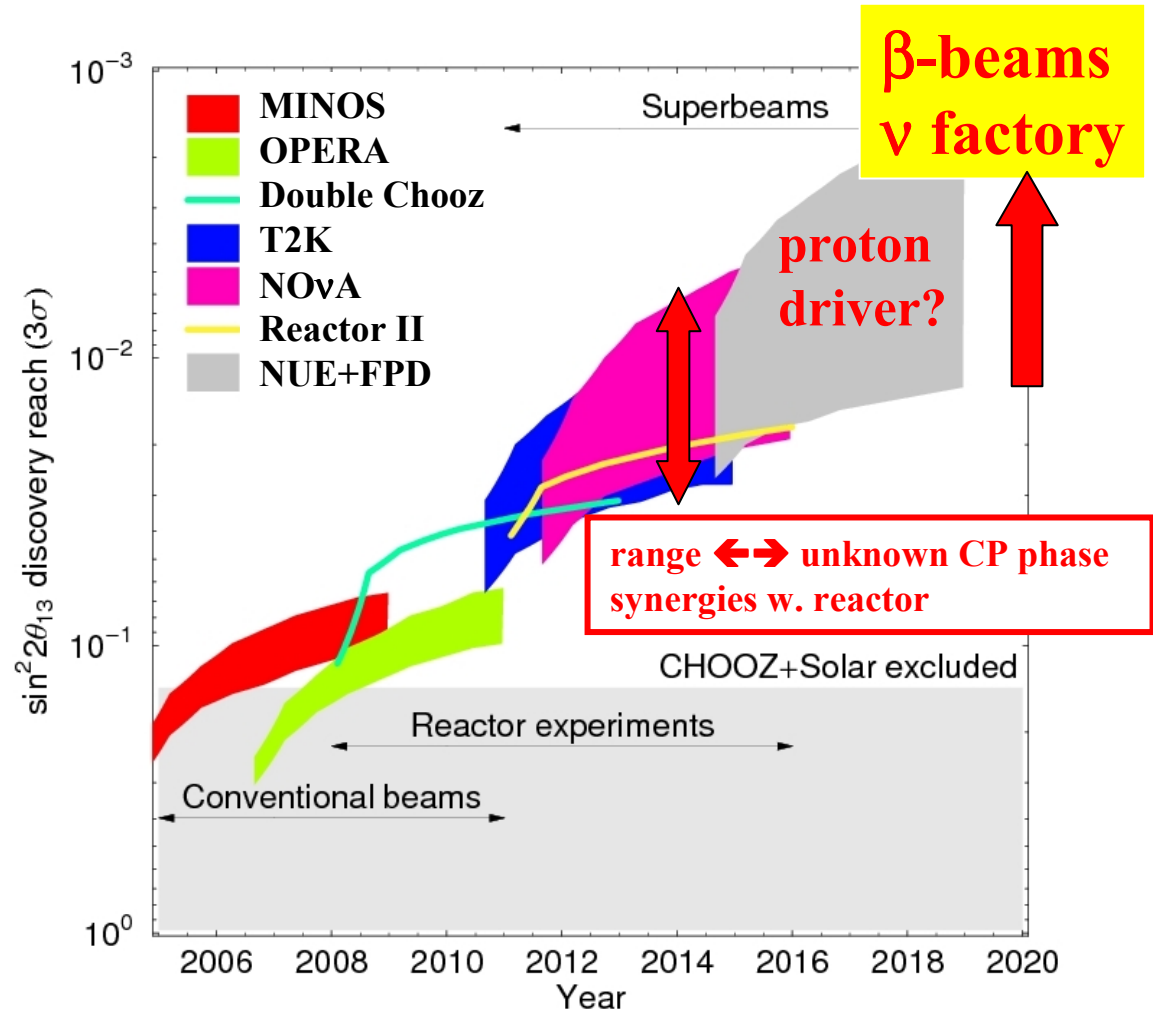
→ correlations & degeneracies, matter effects

θ_{13} – Now and in the Future



Maltoni+Schwetz+Tortola+Valle

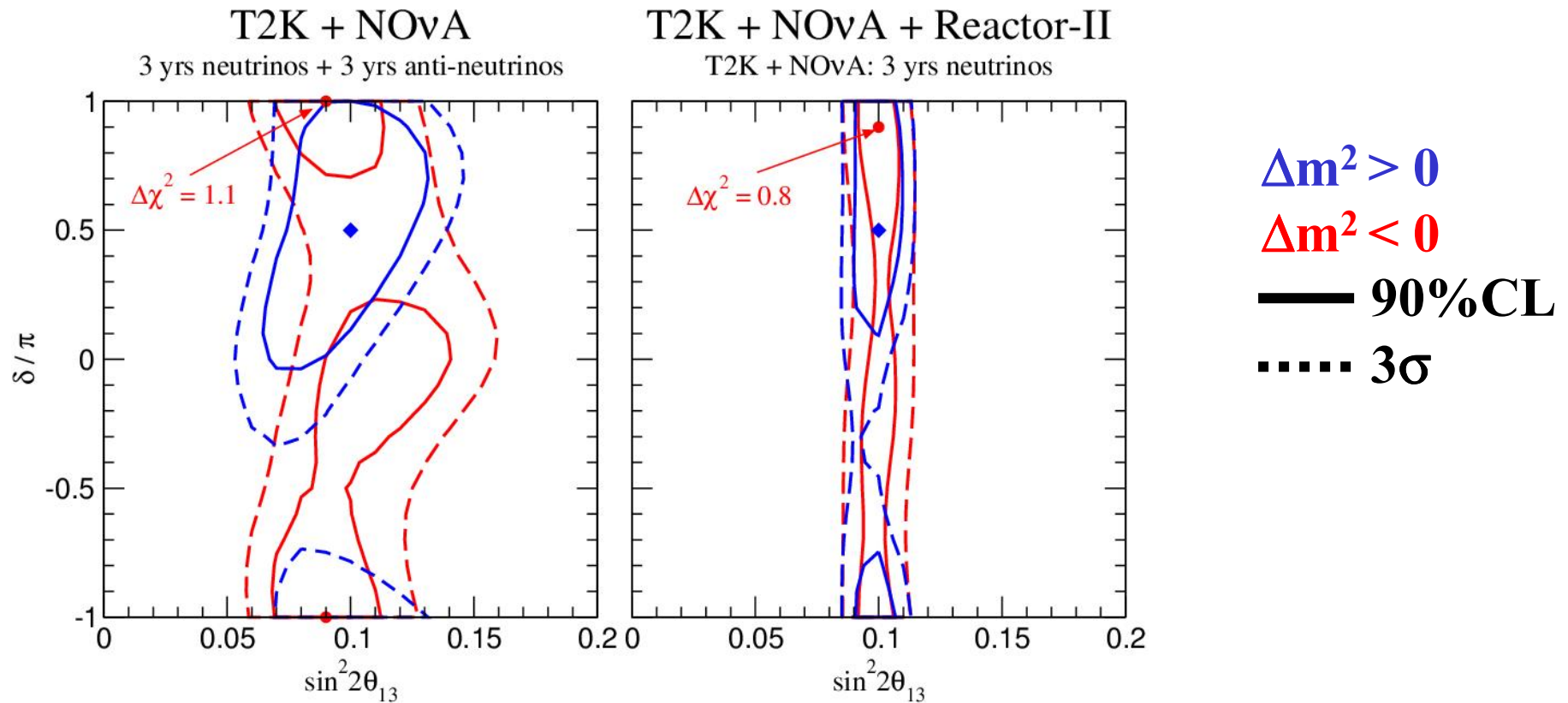
...



Huber, ML, Winter

Leptonic CP-Violation

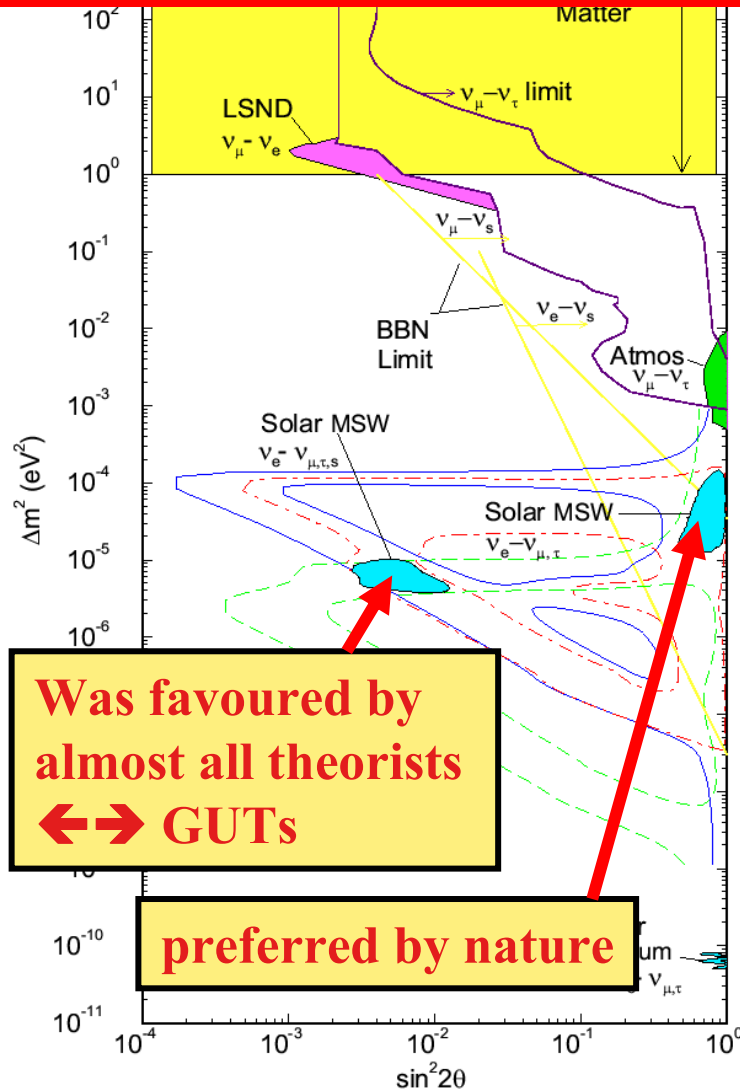
assume: $\sin^2 2\theta_{13} = 0.1$, $\delta = \pi/2$ \rightarrow combine T2K+NOvA+reactor



- \rightarrow bounds or measurements of leptonic CP-violation
- \rightarrow harder for smaller $\sin^2 2\theta_{13}$
- \rightarrow β -beams or/and neutrino factory

Learning about Flavour

History: Elimination of SMA



Next: Smallness of θ_{13} , θ_{23} maximal

- models for masses & mixings
- input: known masses & mixings
 - ➔ distribution of θ_{13} **predictions**
 - ➔ θ_{13} expected close to ex. bound
 - ➔ well motivated experiments

what if θ_{13} is very tiny?
or if θ_{23} is very close to maximal?

- ➔ numerical coincidence unlikely
- ➔ special reasons (symmetry, ...)
- ➔ answered by coming precision

Implications of Precision Measurements

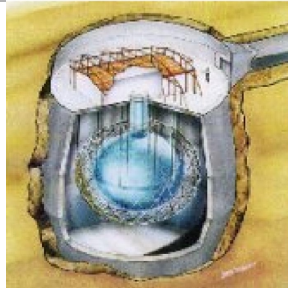
Precision allows to identify / exclude:

- special angles: $\theta_{13} = 0^\circ$, $\theta_{23} = 45^\circ$, ... \leftrightarrow discrete f. symmetries?
- special relations: $\theta_{12} + \theta_C = 45^\circ$? \leftrightarrow quark-lepton relation?
- quantum corrections \leftrightarrow renormalization group evolution

Provides also measurements / tests of:

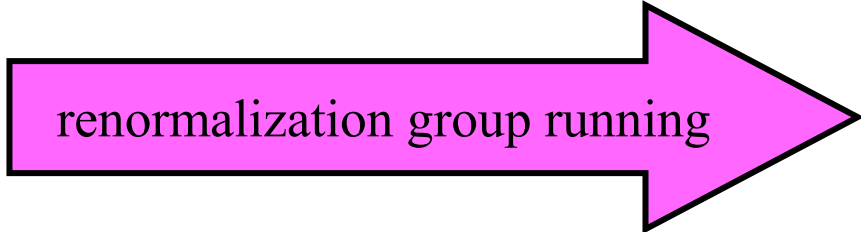
- **MSW effect** (coherent forward scattering and matter profiles)
- **cross sections**
- **3 neutrino unitarity** \leftrightarrow sterile neutrinos with small mixings
- **neutrino decay (admixture...)**
- **decoherence**
- **NSI**
- **MVN, ...**
- \rightarrow **various synergies with LHC and LFV**

Renormalization Group Running



low energies:

- small masses
- large mixings

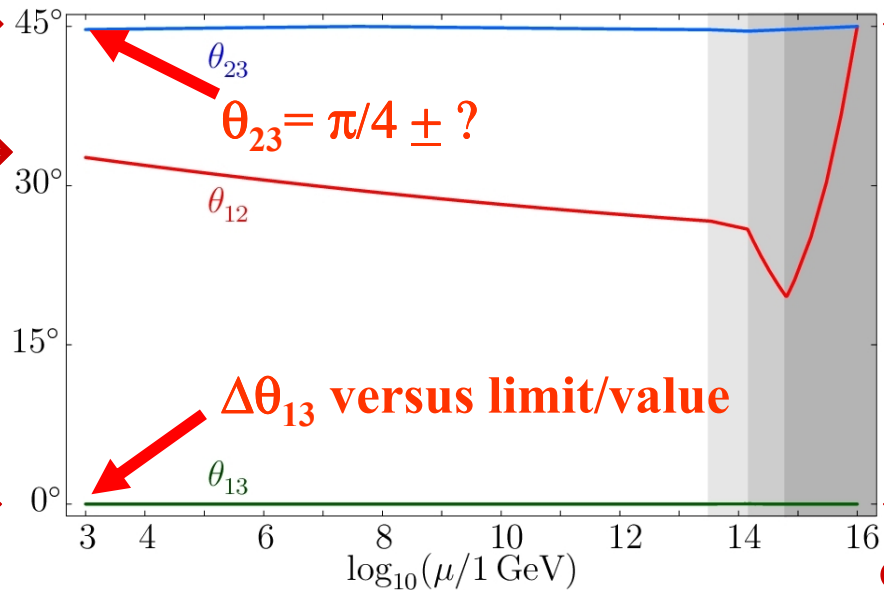


high energies:

- mass models
- flavour-symmetries
- GUT-models, ...

atmospheric \rightarrow 45° \leftarrow bi-maximal

solar \rightarrow



MSSM example:
Antusch, Kersten, ML, Ratz

reactor \rightarrow 0°

\leftarrow Small
or even
zero

The larger Picture: GUTs

Gauge unification suggests that some GUT exists

Requirements:

gauge unification

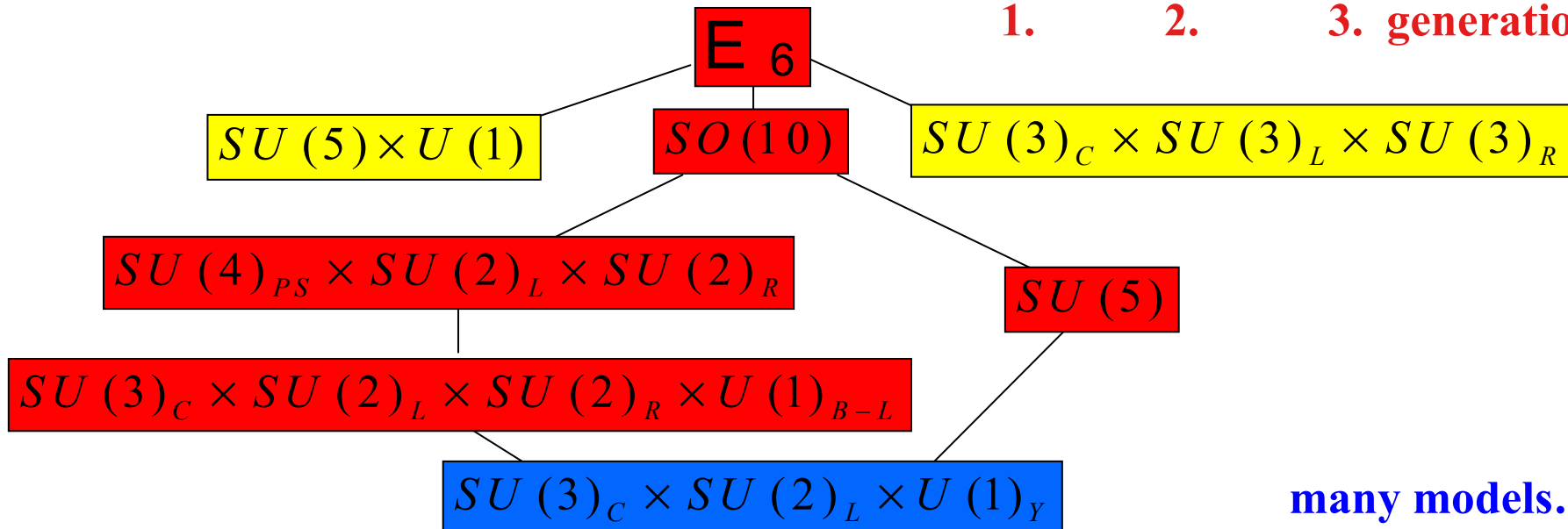
particle multiplets $\leftrightarrow \nu_R$

proton decay

...

Quarks	u	c	t
	d	s	b
Leptons	ν_1	ν_2	ν_3
	e	μ	τ
	$2/3$ ~ 5	$2/3$ ~ 1350	$2/3$ 175000
	$-1/3$ ~ 9	$-1/3$ ~ 175	$-1/3$ ~ 4500
	$0?$	$0?$	$0?$
	0.511	105.66	1777.2

1. 2. 3. generation



many models...

GUT Expectations and Requirements

Quarks and leptons sit in the same multiplets

- one set of Yukawa couplings for given GUT multiplet
- ~ tension: small quark mixings \leftrightarrow large leptonic mixings
- this was in fact the reason for the 'prediction' of small mixing angles (SMA) – ruled out by data

Mechanisms to post-dict large mixings:

- sequential dominance
- type II see-saw
- Dirac screening
- ...

Sequential Dominance

$$m_D = \begin{pmatrix} \cdot & \cdot & \cdot \\ \cdot & a & b \\ \cdot & c & d \end{pmatrix} \quad M_R = \begin{pmatrix} \cdot & \cdot & \cdot \\ \cdot & x & 0 \\ \cdot & 0 & y \end{pmatrix}$$

$$\rightarrow m_\nu = -m_D \cdot M_R^{-1} \cdot m_D^T = \begin{pmatrix} \cdot & \cdot & \cdot \\ \cdot & \frac{a^2}{x} + \frac{b^2}{y} & \frac{ac}{x} + \frac{bd}{y} \\ \cdot & \frac{ac}{x} + \frac{bd}{y} & \frac{c^2}{x} + \frac{d^2}{y} \end{pmatrix}$$

If one right-handed neutrino dominates, e.g. $y \gg x$

\rightarrow small sub-determinant $\sim m_2 \cdot m_3$

$\rightarrow m_2 \ll m_3$ (hierarchy) and $\tan \theta_{23} \simeq a/c$ (large mixing)

$$M_R = \begin{pmatrix} x & 0 & 0 \\ 0 & y & 0 \\ 0 & 0 & z \end{pmatrix} \rightarrow x \ll y \ll z$$

sequential dominance:

$m_1 \ll m_2 \ll m_3$ natural

naturally large mixings

S.F. King, ..

Large/Small versus Maximal/Zero Mixings

- **sequential dominance** → **generically large mixings**
- **experiments will soon tell if**
 - θ_{23} is only large or very close to $\pi/4$
 - θ_{13} is only somewhat small or tiny
- **distinguish mechanism which**
 - produce generically large/small mixings
 - e.g. sequential dominance
 - produce generically only tiny deviations from $\pi/4$ and 0
 - tiny correction to some limiting case
e.g. $\theta_{13} = 0 + \varepsilon$ and $\theta_{23} = \pi/4 - \varepsilon$

Large Mixings and See-Saw Type II

see-saw type II:

- rather natural
- interference of two terms

$$\mathbf{m}_\nu = \mathbf{M}_L - \mathbf{m}_D \mathbf{M}_R^{-1} \mathbf{m}_D^T$$

\mathbf{m}_D and \mathbf{M}_R may have small mixings and hierarchy

However: \mathbf{M}_L can be numerically more important

Example: Break GUT \rightarrow $SU(2)_L \times SU(2)_R \times U(1)_{B-L} \rightarrow \mathbf{M}_L$ from LR

\rightarrow large mixings natural for almost degenerate case $m_1 \sim m_2 \sim m_3$

\rightarrow type I see-saw would only be a correction

type I – type II interference \rightarrow Rodejohann, ML

$\rightarrow \mathbf{M}_L \simeq \mathbf{m}_D \mathbf{M}_R^{-1} \mathbf{m}_D^T \rightarrow$ interesting possibilities

\rightarrow dominance of one term + perturbation by 2nd term

$U_{e3}=0$; maximal θ_{23} \rightarrow small perturbation

Leading structure from one type II term \rightarrow perturbation by 2nd

Three simple, stable candidates for $U_{e3}=0$ and maximal θ_{23}

$$(A) : \sqrt{\frac{\Delta m_A^2}{4}} \begin{pmatrix} 0 & 0 & 0 \\ \cdot & 1 & -1 \\ \cdot & \cdot & 1 \end{pmatrix} \quad L_e \quad EV = \sqrt{\Delta m_A^2} \quad NH$$

$$(B) : \sqrt{\frac{\Delta m_A^2}{2}} \begin{pmatrix} 0 & 1 & 1 \\ \cdot & 0 & 0 \\ \cdot & \cdot & 0 \end{pmatrix} \quad L_e - L_\mu - L_\tau \quad EV = 0 \quad IH$$

$$(C) : m_0 \begin{pmatrix} 1 & 0 & 0 \\ \cdot & 0 & 1 \\ \cdot & \cdot & 0 \end{pmatrix} \quad L_\mu - L_\tau \quad EV = -m_0 \quad \text{degenerate}$$

Perturbation of the Leading Structure

e.g. ‘democratic’ perturbation:

$$m_{\nu}^I \simeq v_L \epsilon \begin{pmatrix} 1 & 1 & 1 \\ \cdot & 1 & 1 \\ \cdot & \cdot & 1 \end{pmatrix}$$

e.g. as correction to case (A):

→ naturally large $\theta_{12} = 1/3$ (tri-bimaximal mixing)

→ finite $\theta_{13} \simeq \sqrt{(\Delta m_{sol}^2 / \Delta m_{atm}^2)} \simeq 1/30$

→ corrections to $\theta_{23} - \pi/4 \simeq \sqrt{(\Delta m_{sol}^2 / \Delta m_{atm}^2)} \simeq 1/30$

Tri-bimaximal Mixing

- **tri-bimaximal mixing works phenomenologically very well**
- **mass matrix can be written as a sum of three terms**

$$m_\nu = \frac{m_1}{6} \begin{pmatrix} 4 & -2 & -2 \\ \cdot & 1 & 1 \\ \cdot & \cdot & 1 \end{pmatrix} + \frac{m_2}{3} \begin{pmatrix} 1 & 1 & 1 \\ \cdot & 1 & 1 \\ \cdot & \cdot & 1 \end{pmatrix} + \frac{m_3}{2} \begin{pmatrix} 0 & 0 & 0 \\ \cdot & 1 & -1 \\ \cdot & \cdot & 1 \end{pmatrix}$$

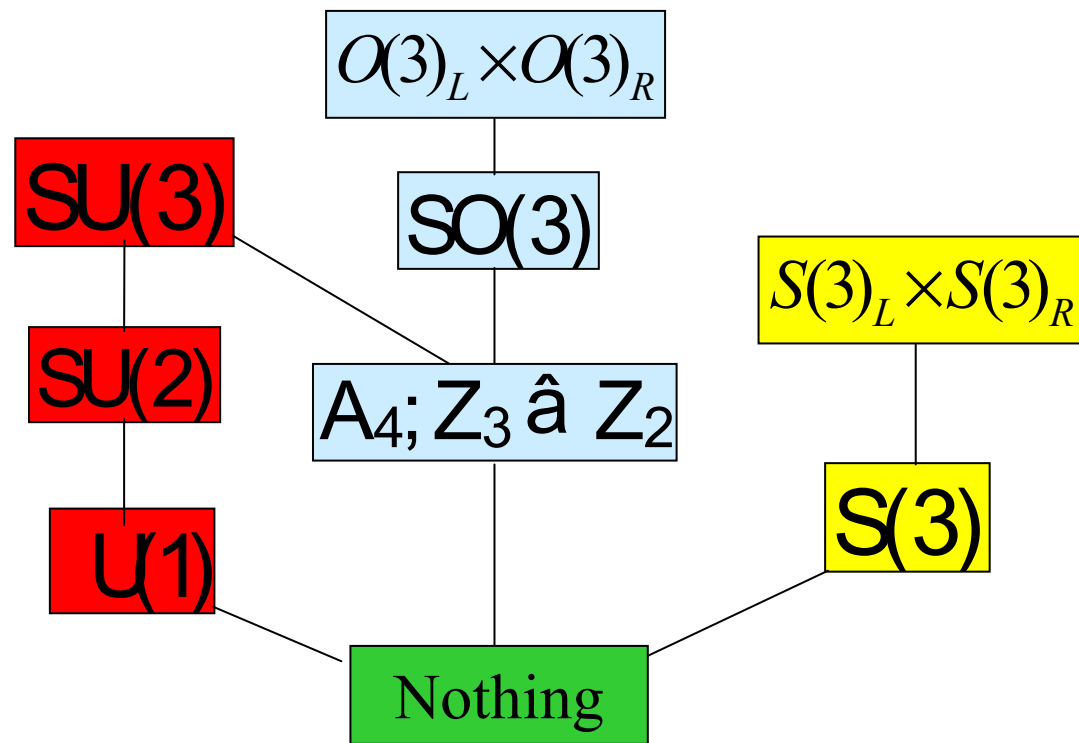
- **phenomenologically very successful**
- **tempting to think of it as a consequence of three terms**
- **type II $\leftrightarrow m_2, m_3$**

Flavour Unification

- so far **no understanding of flavour, 3 generations**
- apparant regularities in quark and lepton parameters
- ➔ flavour symmetries (finite number for limited rank)
- ➔ **symmetry** not texture zeros

Quarks	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
	u ~ 5	c ~ 1350	t 175000
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	d ~ 9	s ~ 175	b ~ 4500
Leptons	0?	0?	0?
	ν_1	ν_2	ν_3
	0.511	105.66	1777.2
	e	μ	τ
	1.	2.	3.
	generation		

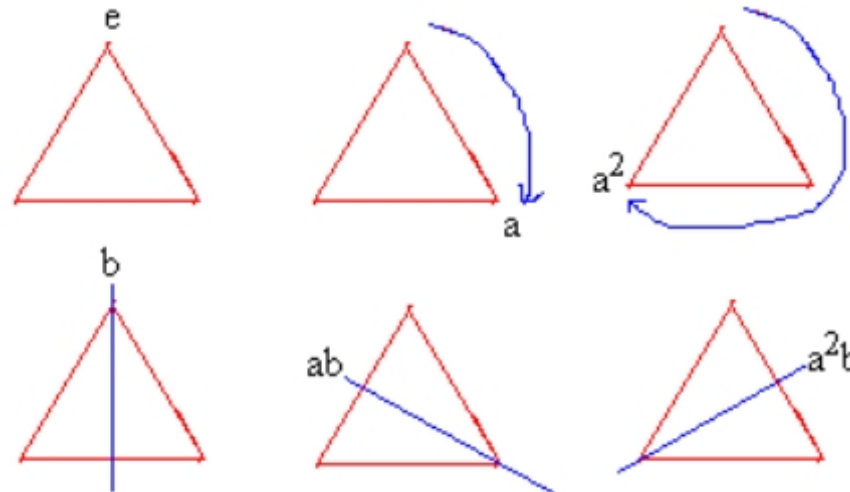
Examples:



Discrete Flavour Symmetries

e.g. dihedral groups D_n

geometric origin of D_3 :



phenomenologically promising example: D_5 Hagedorn, ML, Plentinger

task: search for mass terms which are for suitable Higgs singlets under D_5

1) assign fermions to representations $L = \{L_1, L_2, L_3\}$

2) write down any possible mass term using scalars \leftrightarrow singlet under symmetry

D₅ Allowed Mass Terms

Dirac mass terms: $\lambda_{ij} L_i^T (i\sigma_2) \phi L_j^c$

Majorana mass terms: $\lambda_{ij} L_i^T \equiv \phi L_j$

→ D₅ symmetry induced mass matrices:

Higgses:

$$\Phi_1 \sim 1_1$$

$$\Phi_2 \sim 1_2$$

$$\Psi_1 \sim 2_1$$

L	L^c	Mass Matrix
$(1_2, 1_1, 1_1)$	$(2_1, 1_1)$	$\begin{pmatrix} \kappa_1 \psi_2^1 & -\kappa_1 \psi_1^1 & \kappa_4 \phi^2 \\ \kappa_2 \psi_2^1 & \kappa_2 \psi_1^1 & \kappa_5 \phi^1 \\ \kappa_3 \psi_2^1 & \kappa_3 \psi_1^1 & \kappa_6 \phi^1 \end{pmatrix}$

→ check phenomenology

→ OK + “predictions”

PROBLEM: many successful symmetries

List of models with flavor symmetries (incomplete, by symmetry):

- S_3** : Pakvasa et al. (1978) Derman (1979), Ma (2000), Kubo et al. (2003), Chen et al. (2004), Grimus et al. (2005), Dermisek et al. (2005), Mohapatra et al. (2006), ...
- S_4** : Pakvasa et al. (1979), Derman et al. (1979), Lee et al. (1994), Mohapatra et al. (2004), Ma (2006), Hagedorn, ML and Mohapatra (2006), Caravaglios et al. (2006), ...
- A_4** : Wyler (1979), Ma et al. (2001), Babu et al. (2003), Altarelli et al. (2005,2006), He et al. (2006) ...
- D_4** : Seidl (2003), Grimus et al. (2003,2004), Kobayashi et al. (2005), ...
- D_5** : Ma (2004), Hagedorn et al. (2006).
- D_n** : Chen et al. (2005), Kajiyama et al. (2007), Frampton et al. (1995,1996,2000), Frigerio et al. (2005), Babu et al. (2005), Kubo (2005), ...
- T'** : Frampton et al. (1994,2007), Aranda et al. (1999,2000), Feruglio et al. (2007), Chen and Mahanthappa (2007)
- Δ_n** : Kaplan et al. (1994), Chou et al. (1997), de Medeiros Varzielas et al. (2005), ...
- T_7** : Luhn et al.

GUT \otimes Flavour Unification

$SO(10)$	Quarks	u 2/3 ~5	c 2/3 ~1350	t 2/3 175000
		d -1/3 ~9	s -1/3 ~175	b -1/3 ~4500
	Leptons	ν_1 0?	ν_2 0?	ν_3 0?
		e 0.511	μ 105.66	τ 1777.2
		1.	2.	3.
		$SO(3)_F$		

→ GUT group \otimes flavour group

example: $SO(10) \otimes SU(3)_F$

- SSB of $SU(3)_F$ between Λ_{GUT} and Λ_{Planck}

- all flavour Goldstone Bosons eaten

- discrete sub-groups survive \leftrightarrow SSB

e.g. Z_2, S_3, D_5, A_4

→ structures in flavour space

→ compare with data

GUT \otimes flavour is rather restricted

\leftrightarrow small quark mixings *AND* large leptonic mixings ; quantum numbers

→ so far only a few viable models

Cai and Yu, Hagedorn, ML and Mohapatra, Chen and Mahantappa, King, Ross

→ rather limited number of possibilities; phenomenological success non-trivial

→ aim: distinguish models further by future precision

NSIs & Neutrino Oscillations

Future precision oscillation experiments:

Source	⊗	Oscillation	⊗	Detector
<ul style="list-style-type: none">- neutrino energy E- flux and spectrum- flavour composition- contamination- symmetric $\nu/\bar{\nu}$ operation		<ul style="list-style-type: none">- oscillation channels- realistic baselines- MSW matter profile- degeneracies- correlations		<ul style="list-style-type: none">- effective mass, material- threshold, resolution- particle ID (flavour, charge, event reconstruction, ...)- backgrounds- x-sections (at low E)

precision experiments might see new effects beyond oscillations!

→ modifications of 3f oscillation formulae, different L/E

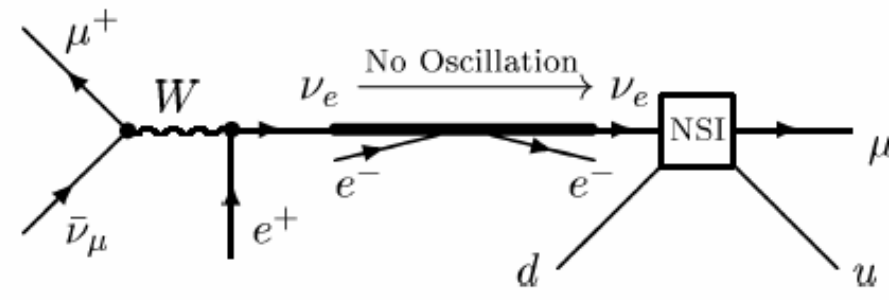
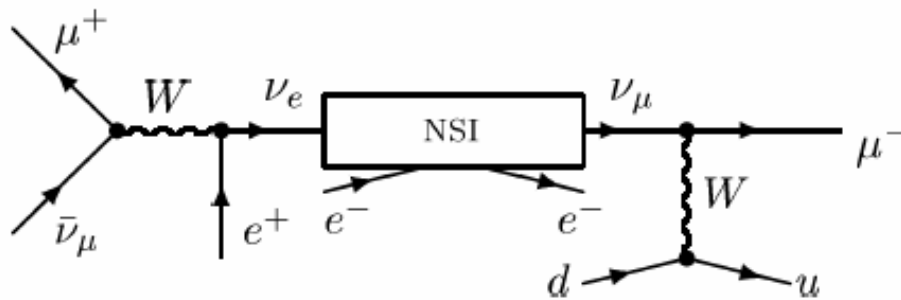
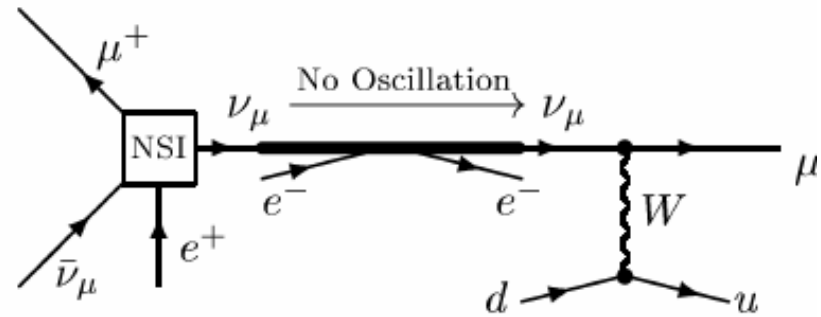
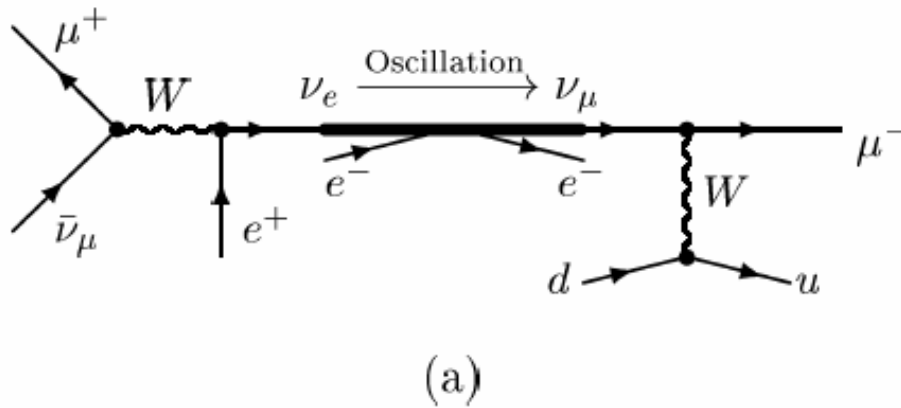
→ small event rates: offset in oscillation parameters

→ Non Standard Interactions = NSI's

NSIs interfere with Oscillations

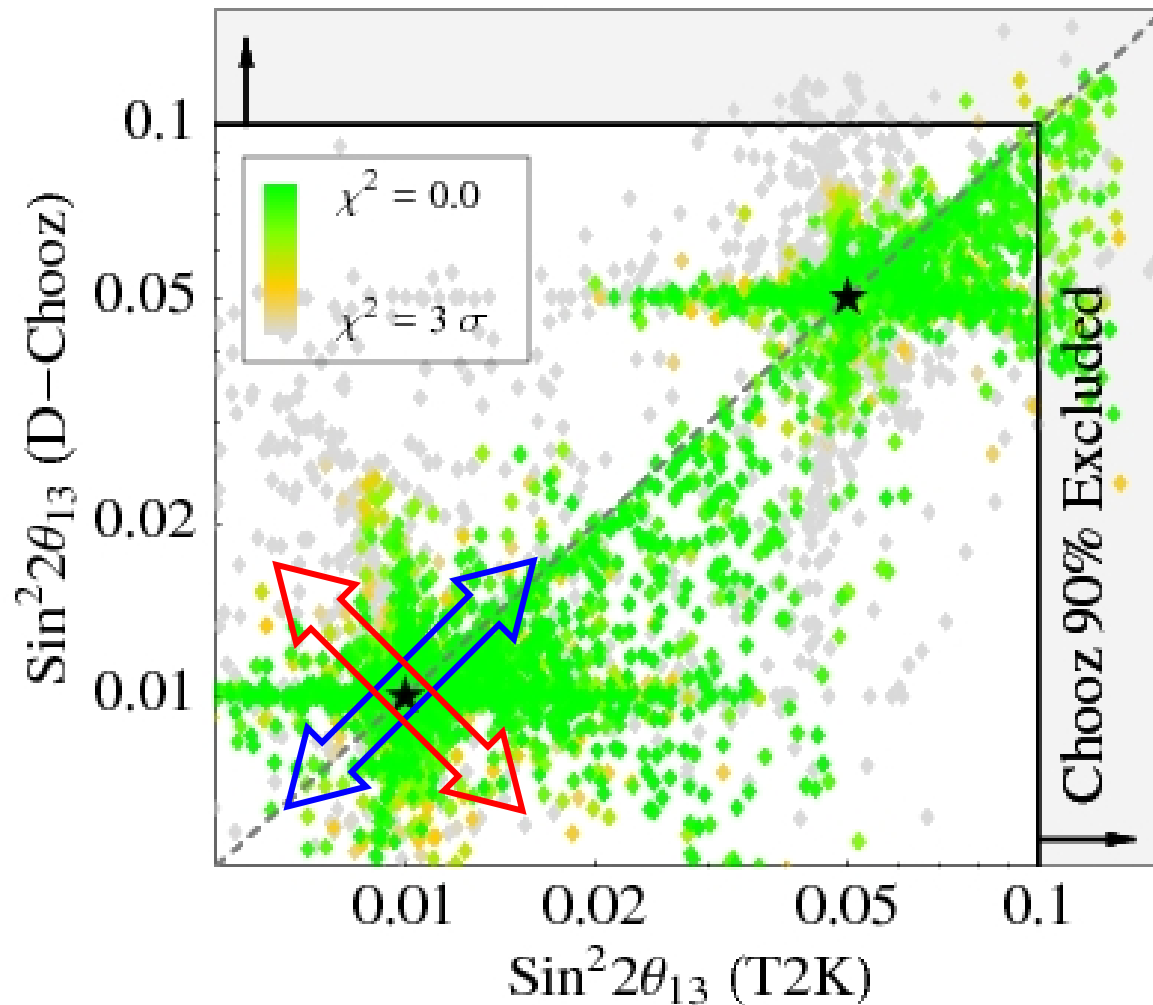
the “golden” oscillation channel

NSI contributions to the “golden” channel



note: interference in oscillations $\sim \epsilon$ \leftrightarrow FCNC effects $\sim \epsilon^2$

NSI: Offset and Mismatch in θ_{13}



Kopp, ML, Ota, Sato

redundant measurement of θ_{13}

Double Chooz + T2K

***=assumed 'true' values of θ_{13}**

scatter-plot:

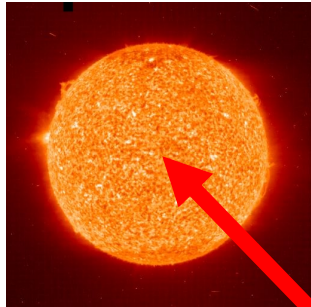
- ϵ values random
- below existing bounds
- random phases

NSIs can lead to:

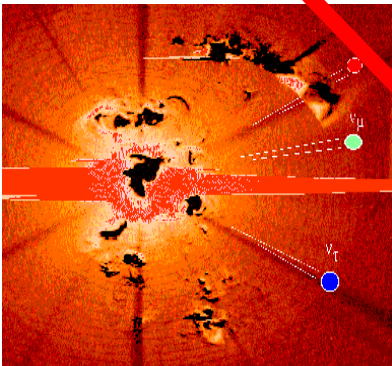
- **offset**
- **mismatch**

- ➔ **redundancy**
- ➔ **interesting potential**

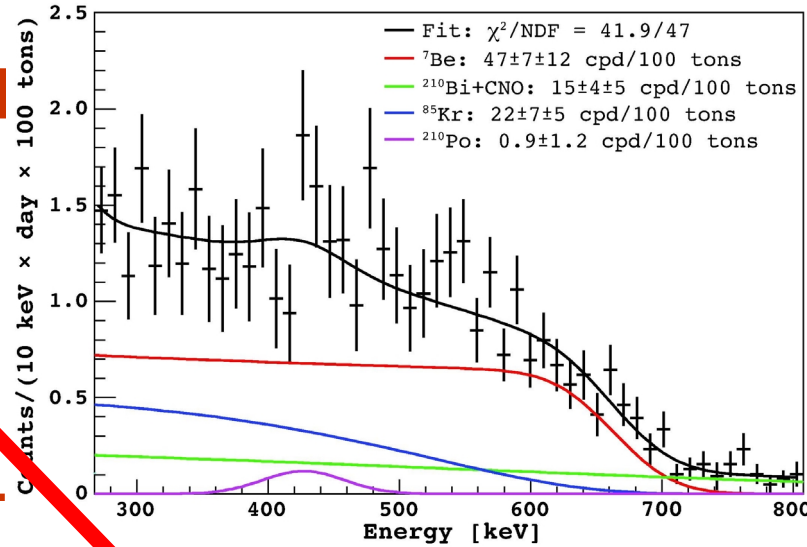
New Physics & Neutrino Sources



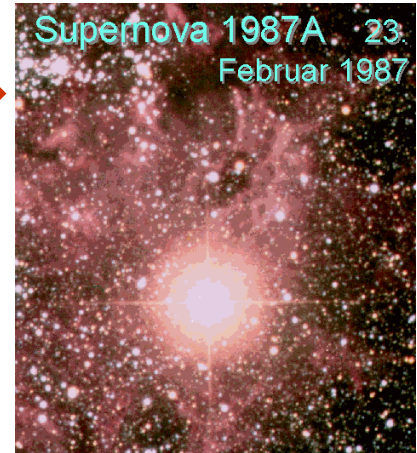
← **Su**



←



y: →
ae



→ **accelerators**

BOREXINO

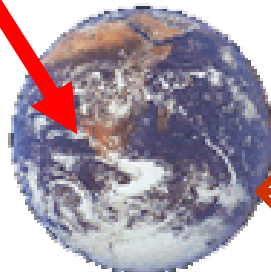
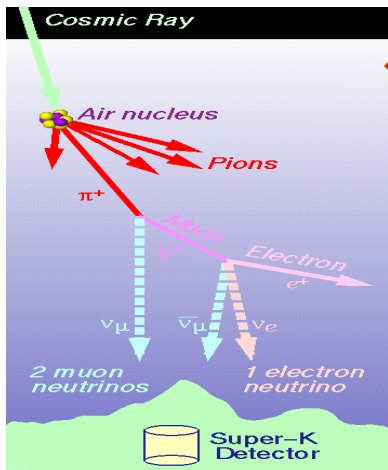
more to come:

← **Atmosphere**

→ **3.5% seasonal variation...**

→ **geo-neutrinos, ...**

Accelerators →



← **Earth**



Conclusions

- **neutrino physics very promising → unique information**
 - sources → important connections to various fields
 - particle physics properties
- **first solid particle physics beyond the standard model**
 - up to 9 more parameters
 - 1st explicit fermion mass terms
 - lepton number violation violation
- **future: precision neutrino physics**
 - very precise measurements
 - NSIs: offsets & mismatch possible → redundancy
 - synergies with LHC and LFV physics
- **interpreting & understanding flavour structures**
 - flavour symmetries - GUTs - SUSY
 - ...various fancier ideas

Neutrinos probe new physics in many ways!

