

TAUP 2007

11-15 September 2007, Sendai, Japan

The ArDM
a ton-scale liquid argon experiment
for direct
dark matter detection

A. Badertscher, L. Kaufmann, L. Knecht,
M. Laffranchi, P. Lightfoot, A. Marchionni,
G. Natterer, P. Otyugova, A. Rubbia, J. Ulbricht
ETH Zurich, Switzerland

C. AMSLER, V. Boccone,
S. Horikawa, C. Regenfus, J. Rochet
Zurich University, Switzerland

A. Bueno, M.C. Carmona-Benitez, J. Lozano,
A. J. Melgarejo, S. Navas-Concha, A. Ruiz
University of Granada, Spain

M. Daniel, P. Ladron de Guevara, L. Romero
CIEMAT, Spain

P. Mijakowski, P. Przewlocki, E. Rondio
Soltan Institute Warszawa, Poland

H. Chagani, P. Majewski,
M. Robinson, N. Spooner
University of Sheffield, England

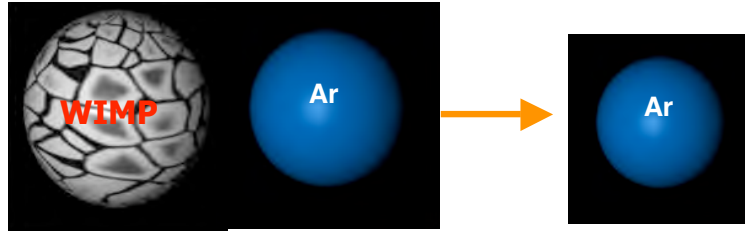
P.Otyugova, ETH Zurich

ArDM WIMP detection mechanism

WIMP-Argon elastic scattering

A WIMP collides with argon inside the detector...

...transmitting its kinetic energy to the nucleus...



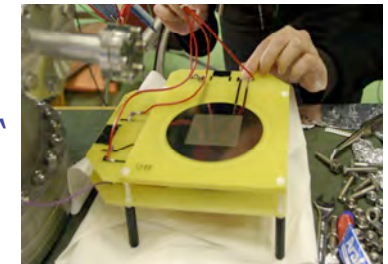
ArDM is a one ton liquid argon detector designed to measure ionization charge with a good spacial resolution and scintillation light.

The light is „seen“ by photomultipliers



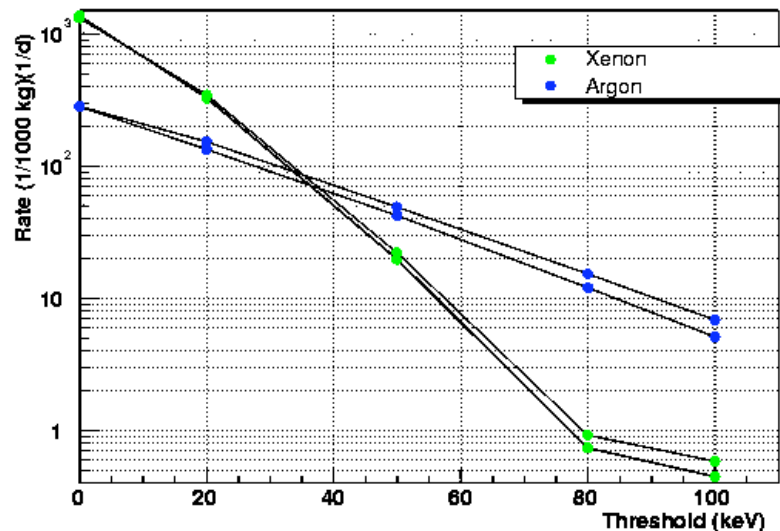
Light and free electrons are produced from interaction with neighbouring argon atoms

The electrons are „seen“ by electron multipliers

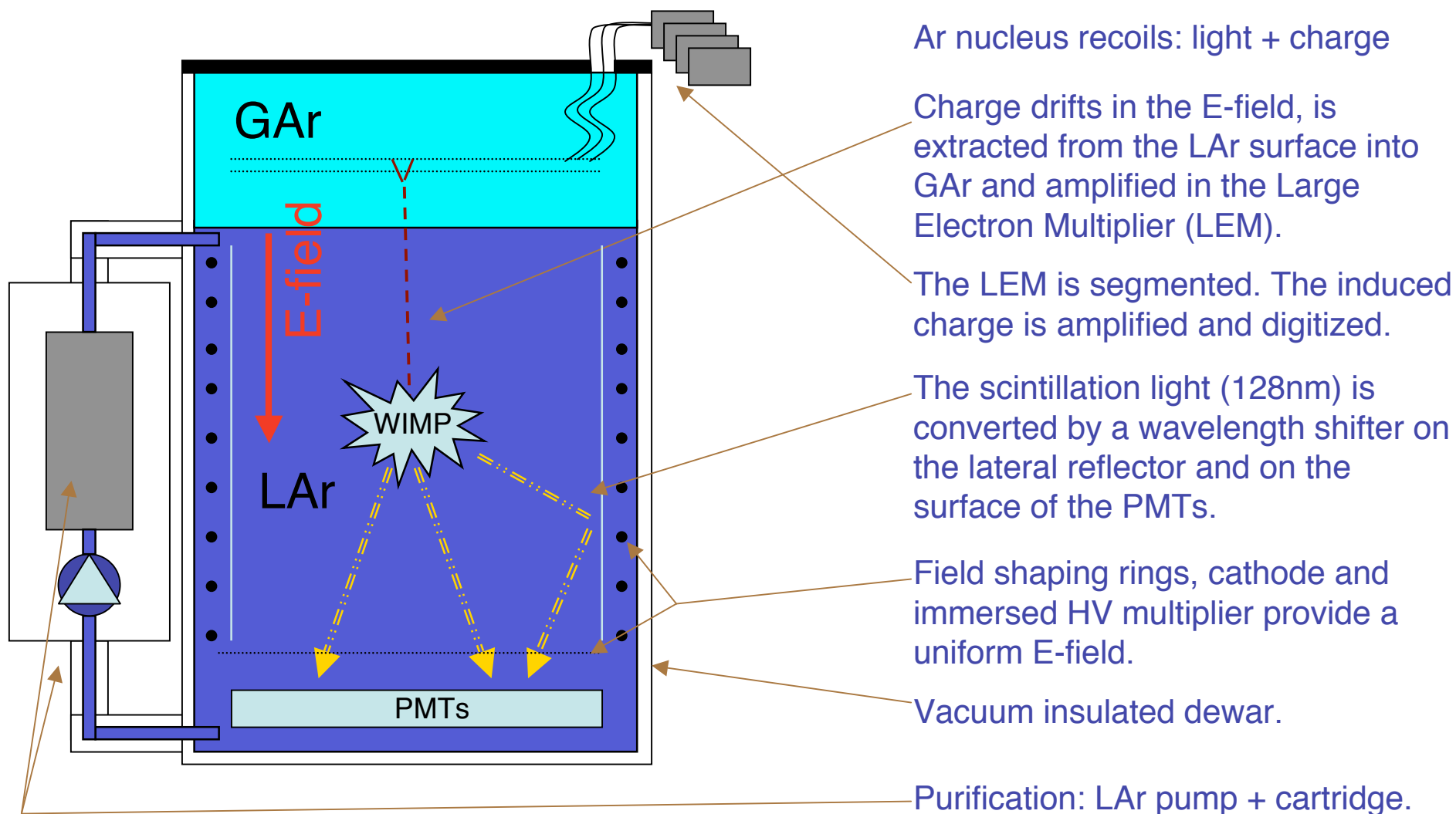


Assumptions for simulation:

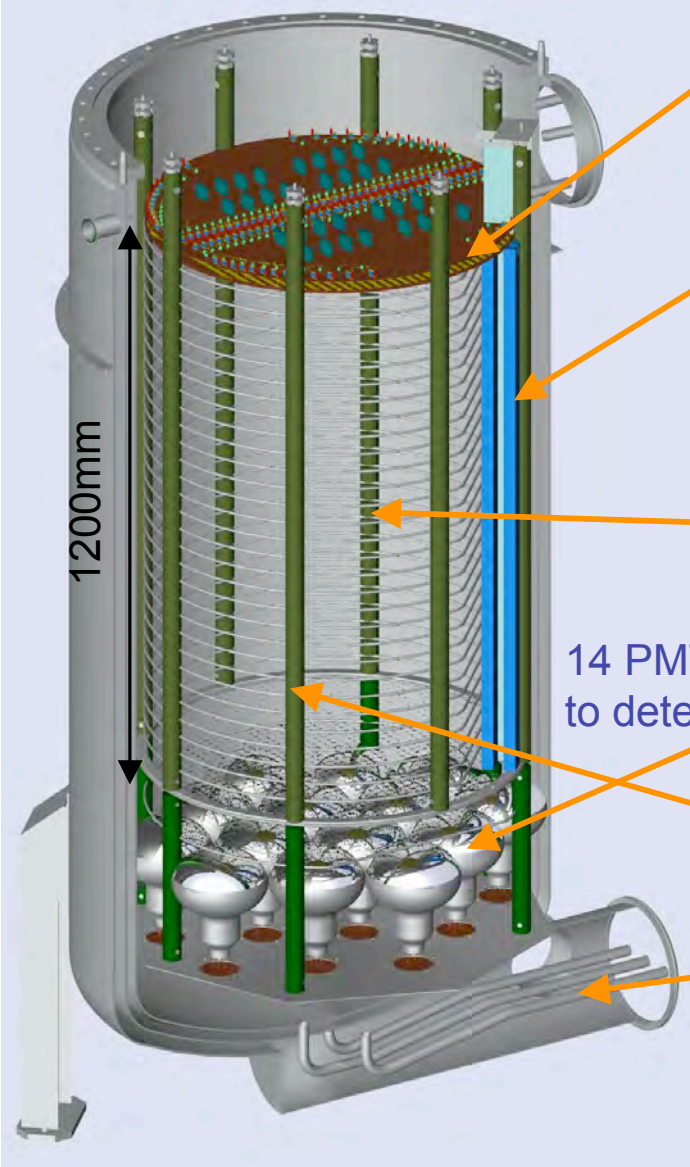
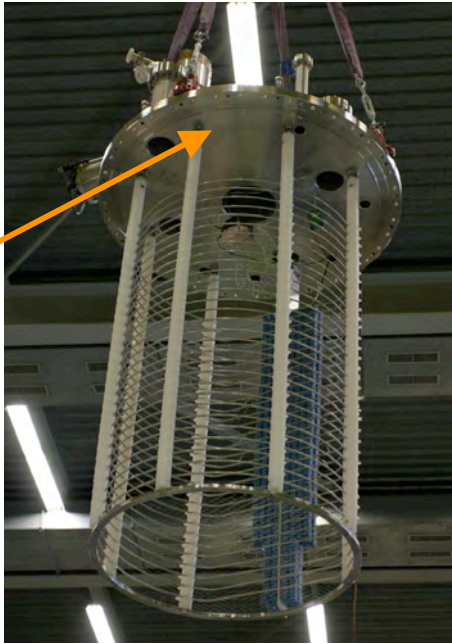
- **Cross-section normalized to nucleon**
 - $\sigma = 10^{-42} \text{ cm}^2 = 10^{-6} \text{ pb}$
 - $M_{\text{WIMP}} = 100 \text{ GeV}$
- **Halo Model**
 - WIMP Density = 0.5 GeV/cm^3
 - $v_{\text{esc}} = 600 \text{ km/s}$
- **Interaction**
 - Spin independent
 - Engel Form factor



ArDM bi-phase detection principle



Detector Layout



Two-stage LEM.
800mm diameter

Greinacher chain:
supplies the right
voltages to
the field shaping rings
and the cathode

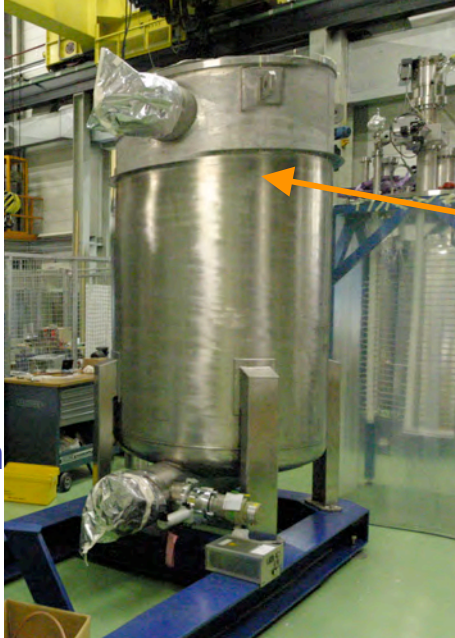
Detector inner
part with the upper
flange

Field shaping
rings

14 PMTs below the cathode
to detect the scintillation light.

Support
pillars

Input/output
of the recirculation
system



ArDM
Dewar

Summary of the detector parameters

Detector	
Max. drift length	120 cm
Target mass	850 kg
High voltage	
Drift field	1-5 kV/cm
Charge readout	
LEM gain	10^4 per e^-
Light readout	
Global collection efficiency	3%

Charge Read-Out System: Large Electron Multiplier (LEM)

GEM: Ref. F.Sauli, NIM A, 1997, vol. 386, p.351

THGEM: Ref. Chechik.R., Breskin,A., Shalem,C.,Mormann,D.,
NIM A, 2004, vol. 535, p.303

LEM is a thick macroscopic GEM

Diameter of the hole: 500 microns.
Distance between two holes: 800 microns.

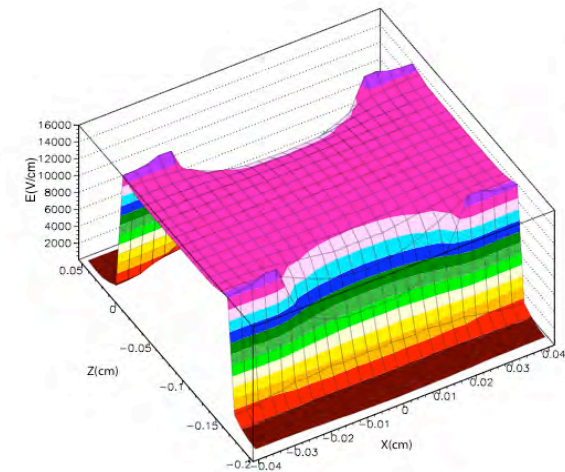
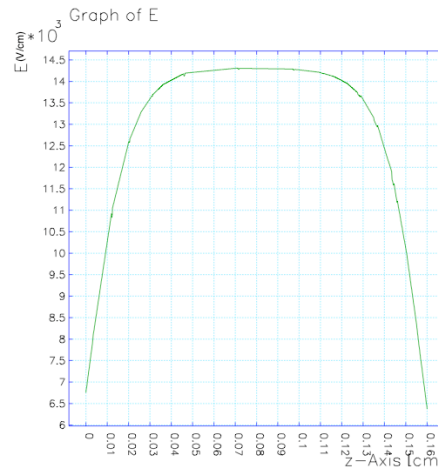
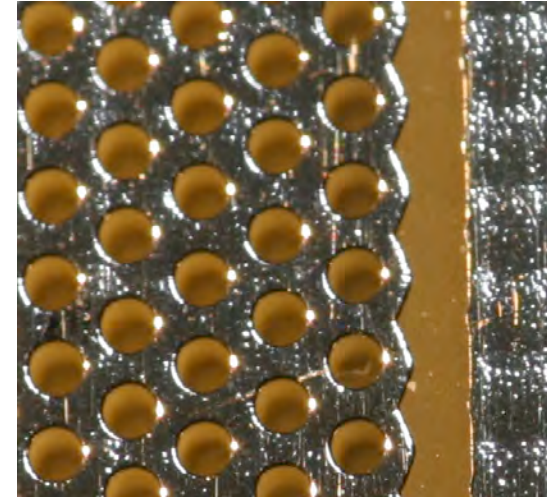
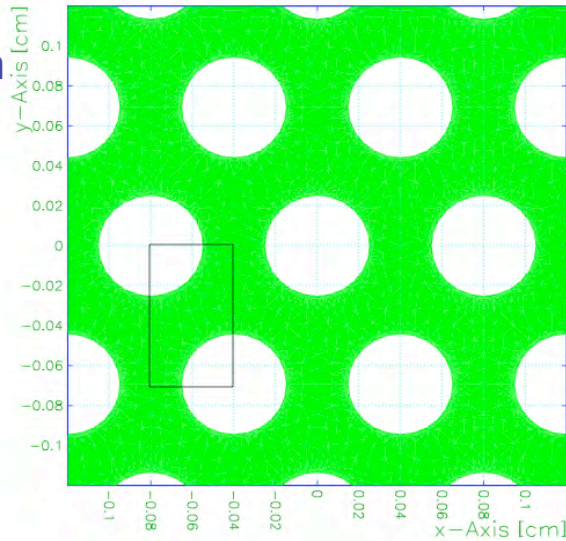
LEM thickness: 1.5mm.

For HV supply both surfaces are covered with copper electrodes.

LEM is manufactured on standard PCB technique. The holes are produced by drilling. Copper electrodes are covered with palladium layer in order to avoid oxidization.

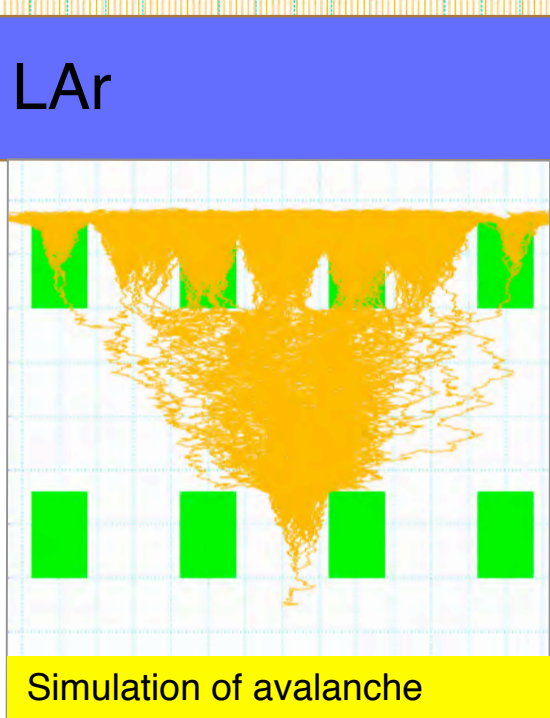
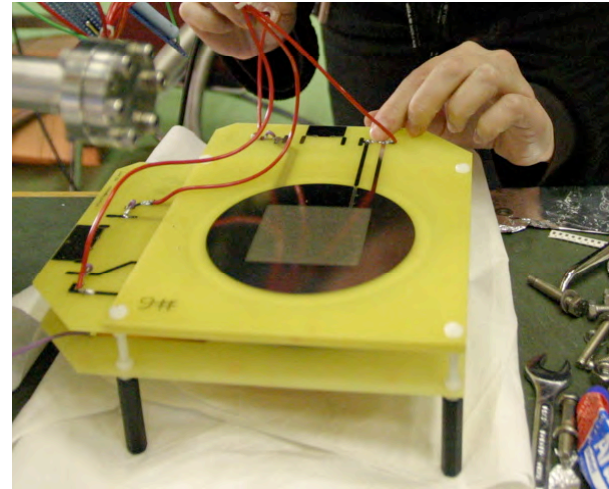
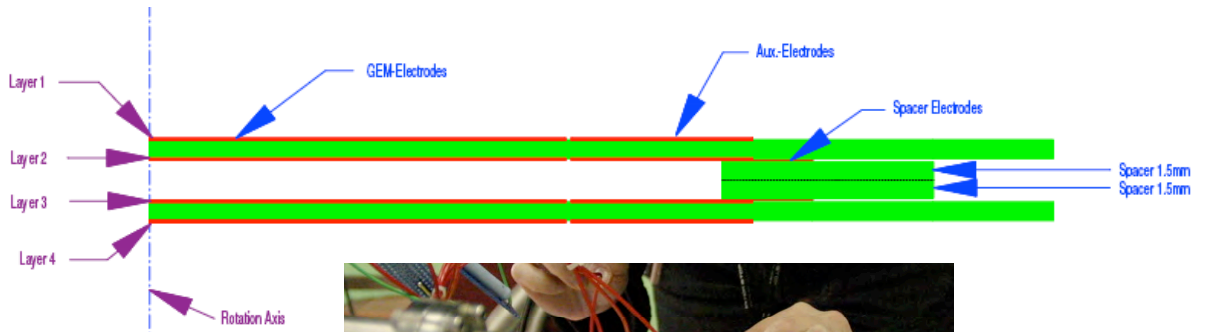
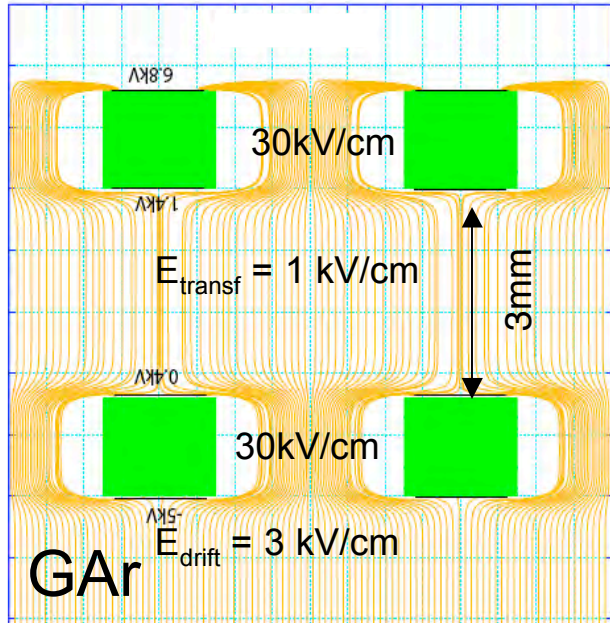
Thickness of the electrodes is 35 microns.

Layout of the cell

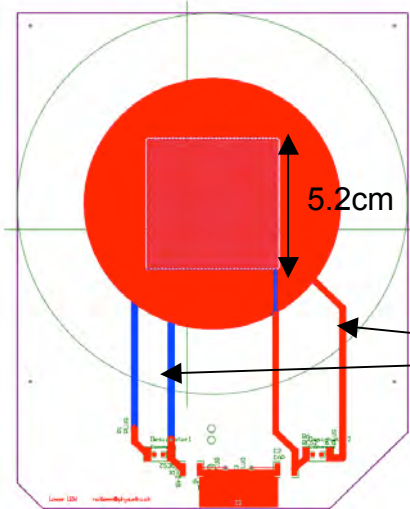


P.Otyugova, ETH Zurich

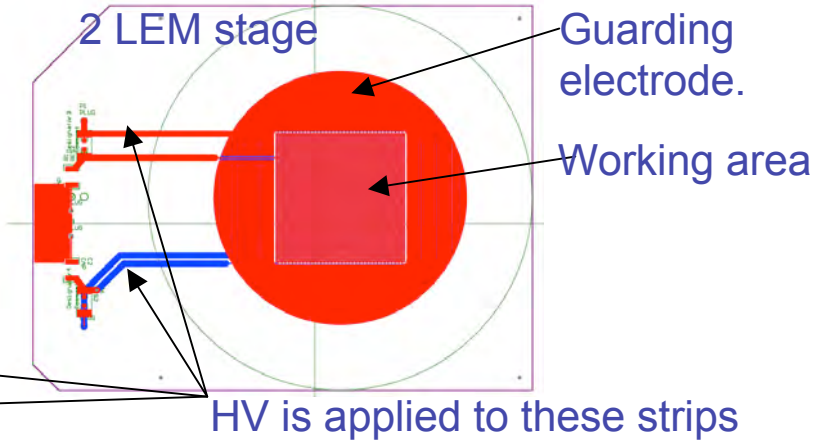
Double-stage LEM system



1 LEM stage

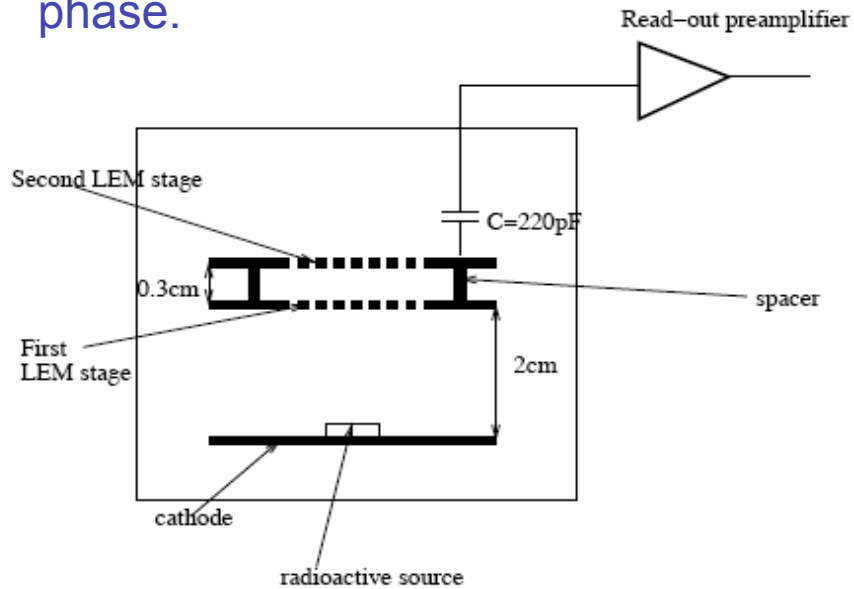


2 LEM stage

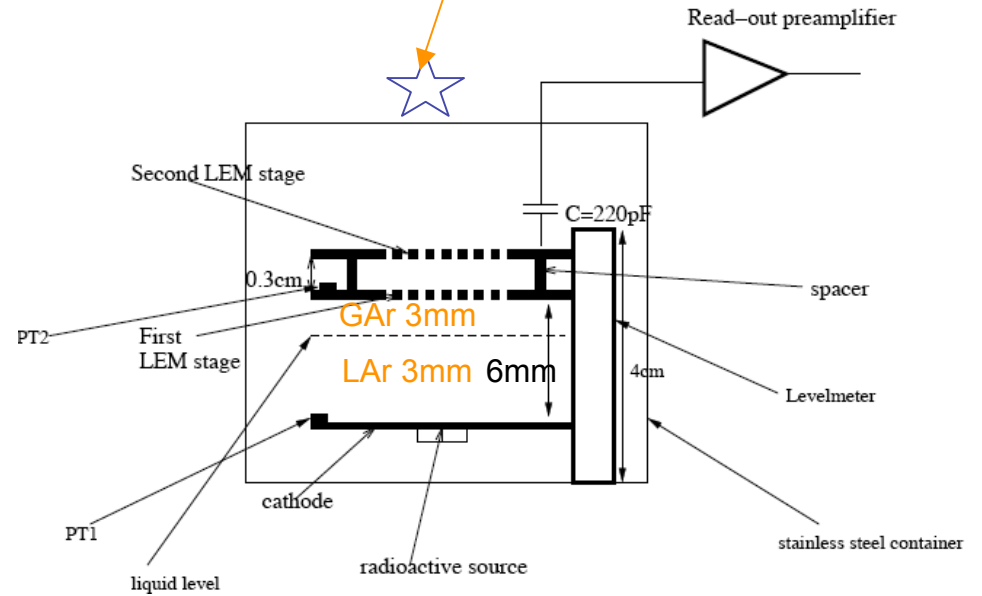


Experimental setups

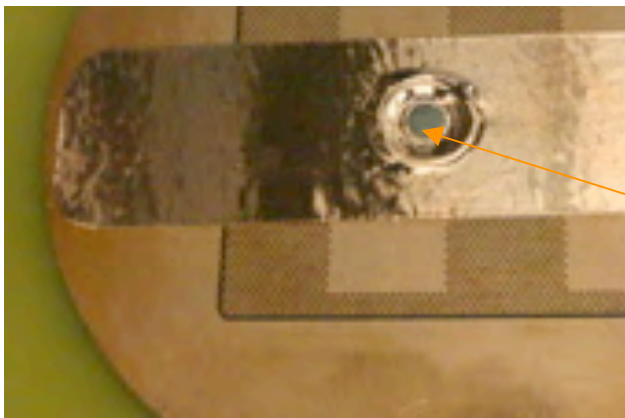
Setup for measurements in single gas phase.



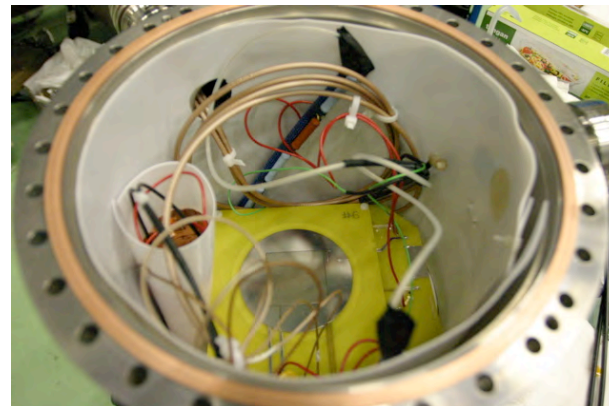
External radioactive sources,
 Cs^{137} 662keV 240kBq.
 Co^{60} 1.17,1.33MeV 4.85kBq



Setup for measurements in double phase

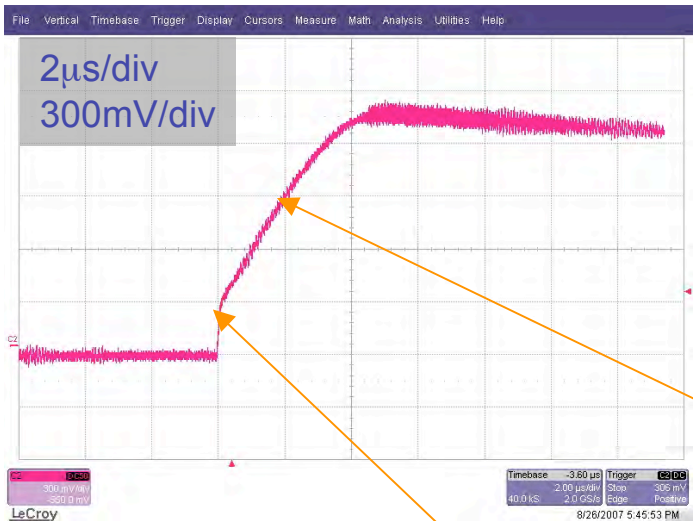


Internal r/a source
 Fe^{55} , 5.9keV,
 12kBq



Signal shapes

Signals have different shapes in pure Ar and in 90% Ar 10% CO₂ mixture. These signals were measured at room temperature and at atmospheric pressure.

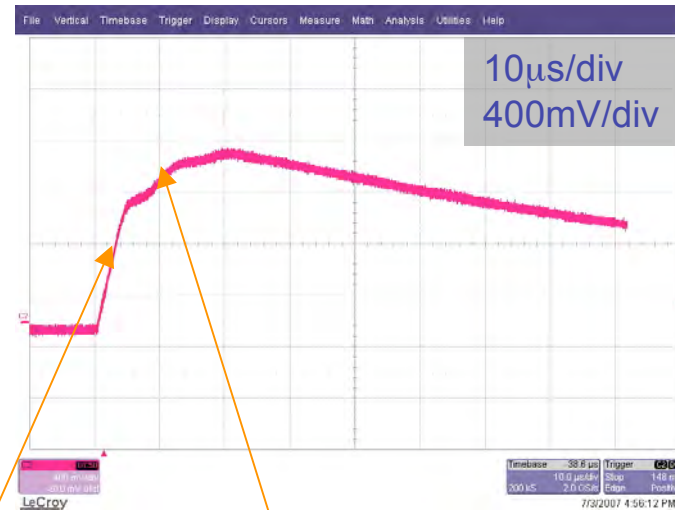


Ion- induced signal (5 μs)

Electron- induced O(100ns) signal

Signal shape in ArCO₂ mixture
risetime is about 5 μs.

The signal risetime has only one ion-induced component.
photons are absorbed by CO₂.

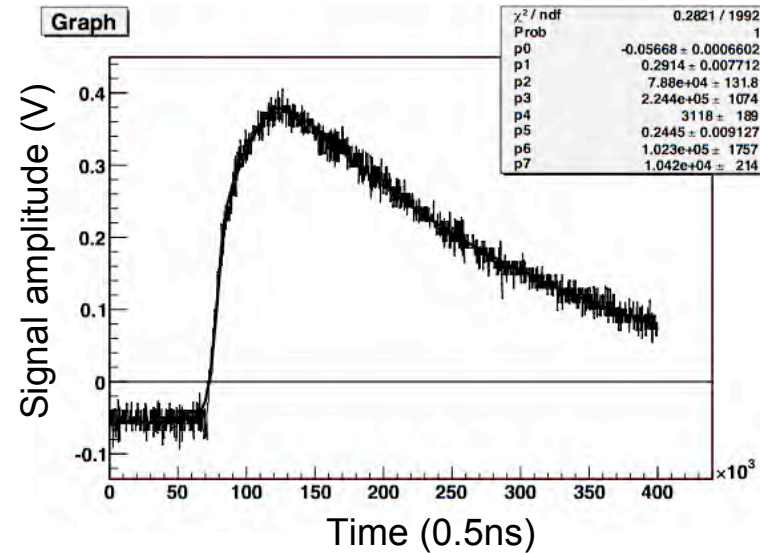
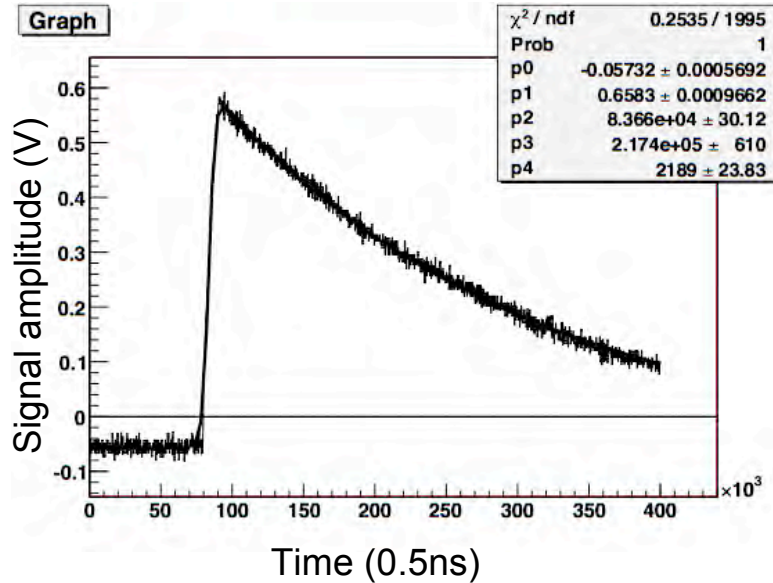


Signal shape in pure Ar.
Risetime is about 20 μs.
2 components of the risetime are visible.

Fast ion-induced component, coming from development of a primary avalanche (5 μs)

Slow ion-induced components, coming from development of a secondary photo-avalanche. (O(15-20 μs)).

Signal fit functions.



$$f(t) = B + A \frac{e^{-\frac{(t-t_0)}{\tau_1}}}{1 + e^{-\frac{(t-t_0)}{\tau_2}}}$$

$$f_1(t) = B + A \frac{e^{-\frac{(t-t_0)}{\tau_1}}}{1 + e^{-\frac{(t-t_0)}{\tau_2}}} + C \frac{e^{-\frac{(t-t'_0)}{\tau_1}}}{1 + e^{-\frac{(t-t'_0)}{\tau'_2}}}$$

B -baseline;

t_0 -point for which the height of the function with respect to the baseline is equal to $A/2$;

A -related to the amplitude of the fast component;

τ_1 τ_2 -are related to risetime of a fast component and a falltime respectively;

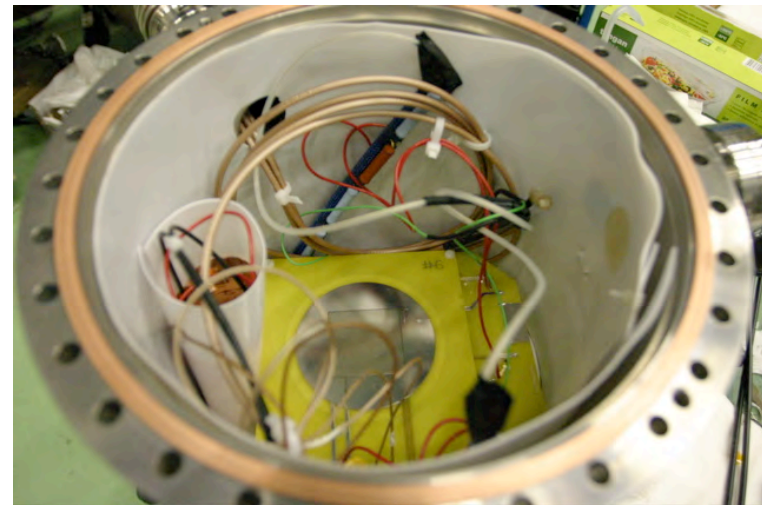
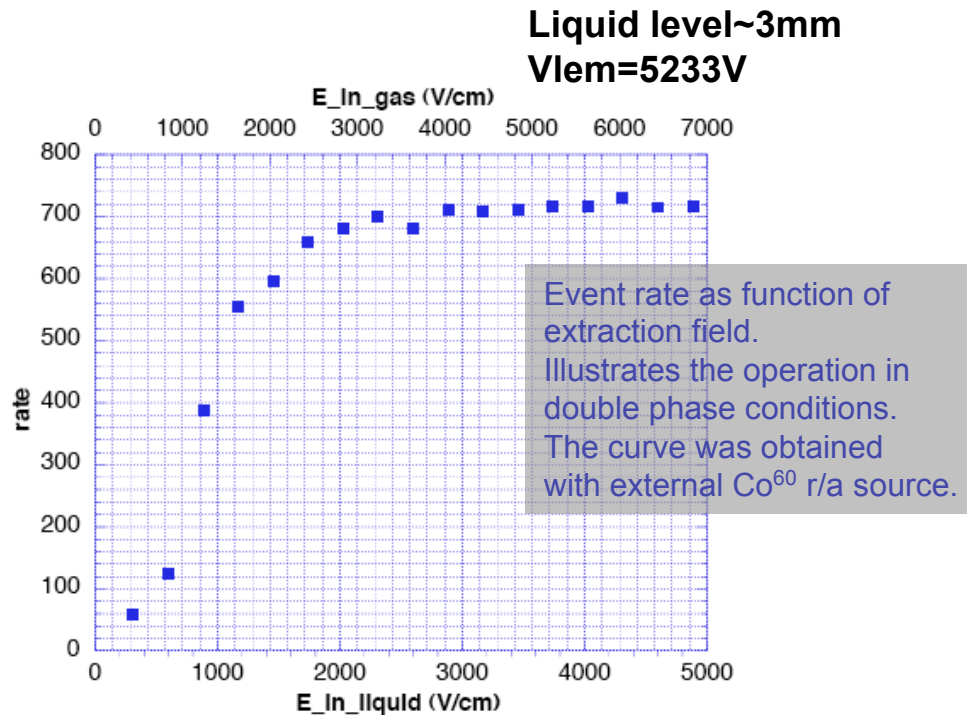
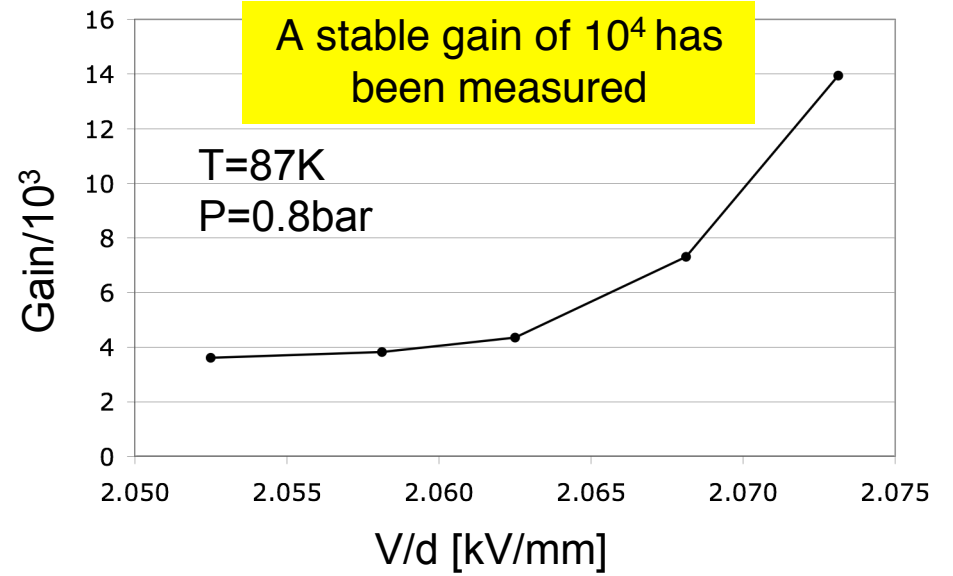
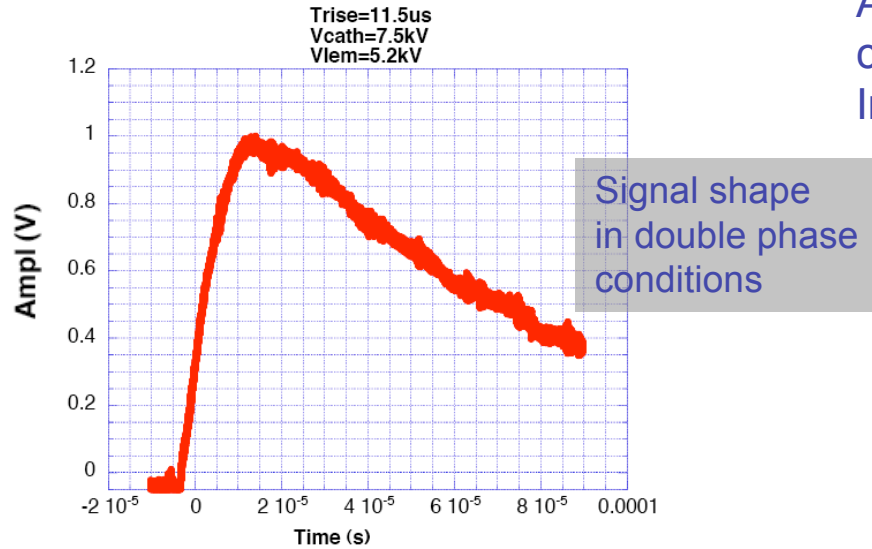
C -related to the amplitude of the slow component;

t'_0 -point for which the height of the function with respect to the baseline is equal to $C/2$;

τ'_2 -is related to the risetime of a slow component.

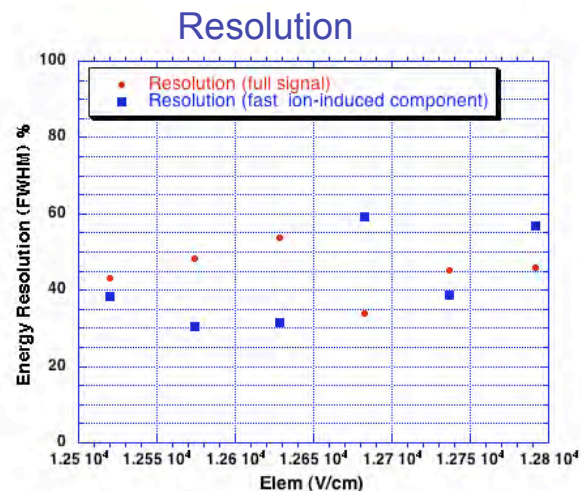
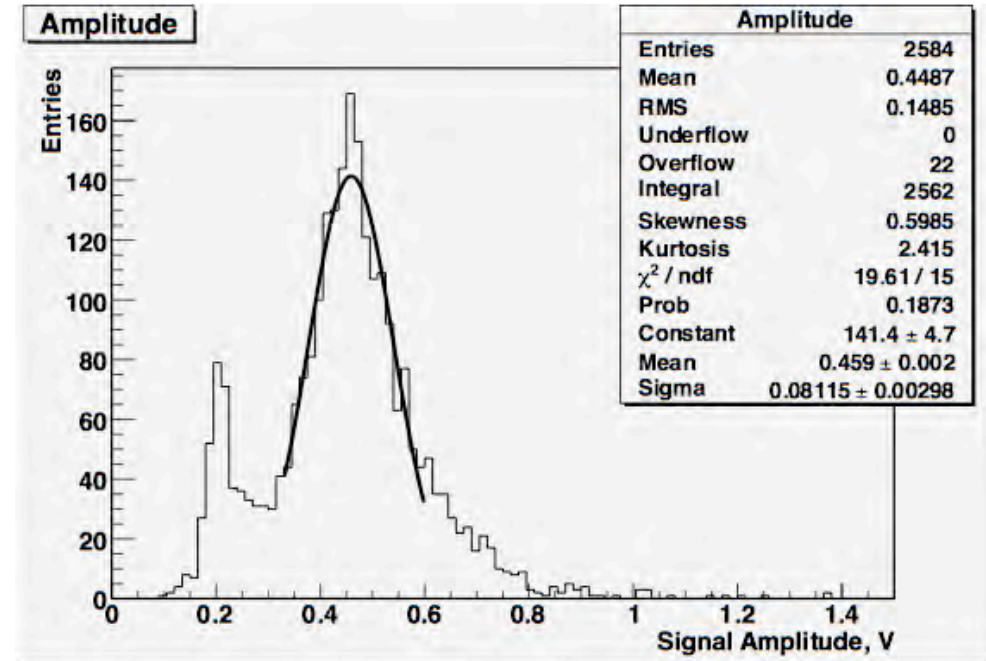
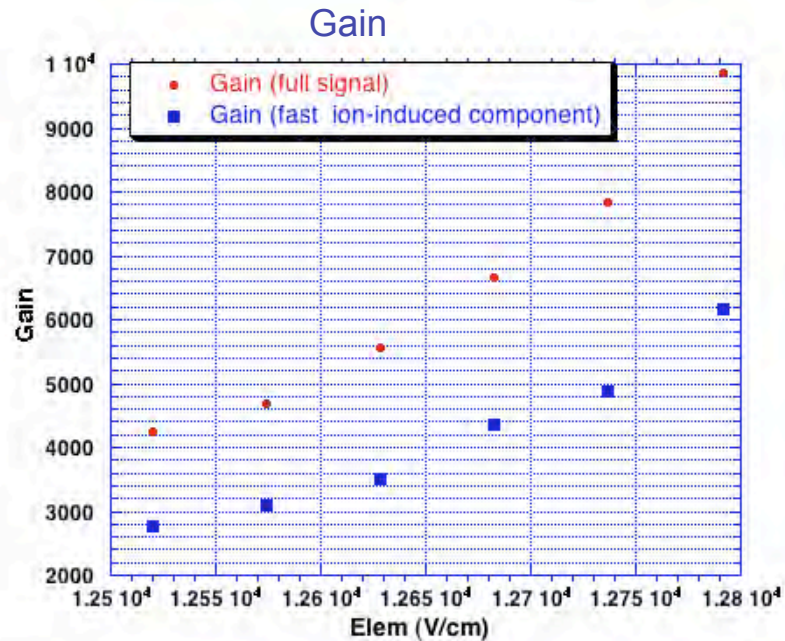
Tests at cryogenic temperatures and double phase conditions.

A stable gain of 10^4 was obtained in the gas phase at cryogenic temperatures. Curve was obtained with Fe^{55} Internal r/a source.



Gain estimation and signal amplitude distribution.

Resolution (FWHM)=42.5%



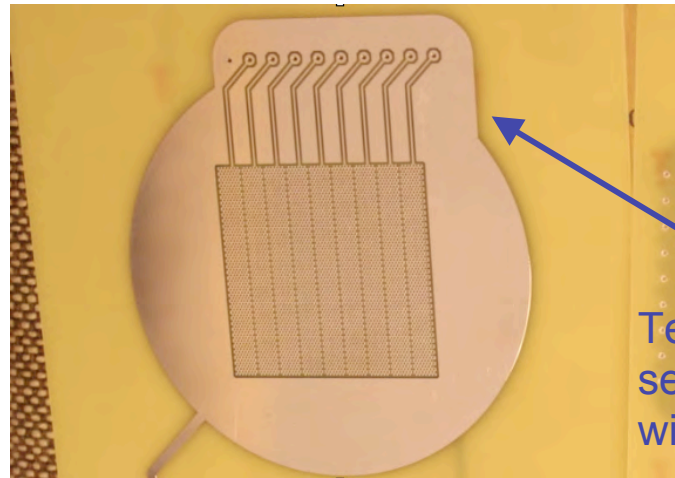
Conditions:

$V_{\text{lem}} = 1.9\text{kV}$
 $V_{\text{cath}} = 2.5\text{kV}$
 Electric field:
 $E = 12.6\text{kV/cm}$
 Drift field:
 $E_d = 0.53\text{kV/cm}$

Amplitude distribution was obtained with pure Ar gas at atmospheric pressure and room temperature. R/a source: Fe^{55} , 5.9keV. The source was collimated to the diameter of 1mm in order to decrease the event rate.

Source Rate: 240Hz.

Segmented LEM

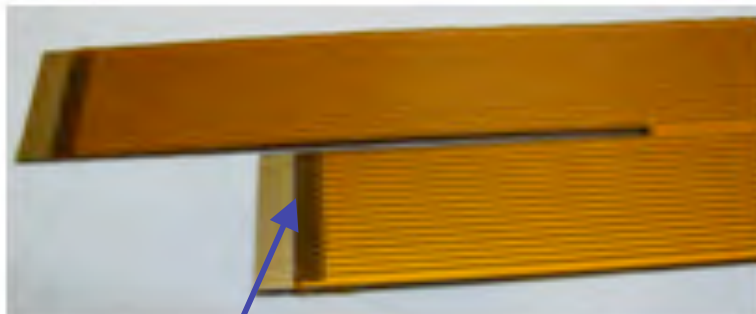


Test prototype of a segmented LEM. Strip width: **6mm**

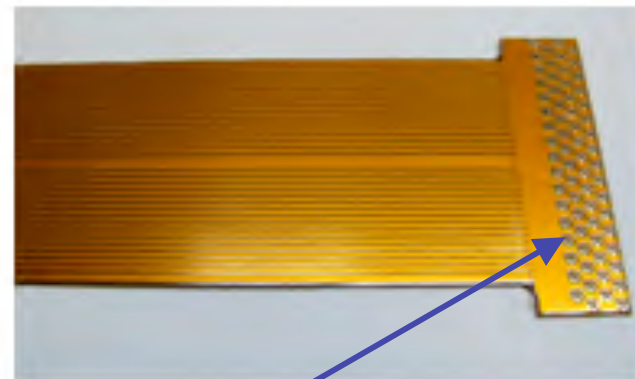
Final LEM charge readout system will be segmented.

Electrodes on both sides are striped. Strips are perpendicular to each other.

Final number of channels: **1024**
Strip width: **1.5mm**



to ZIF connector on the LEM board



To the preamplifiers

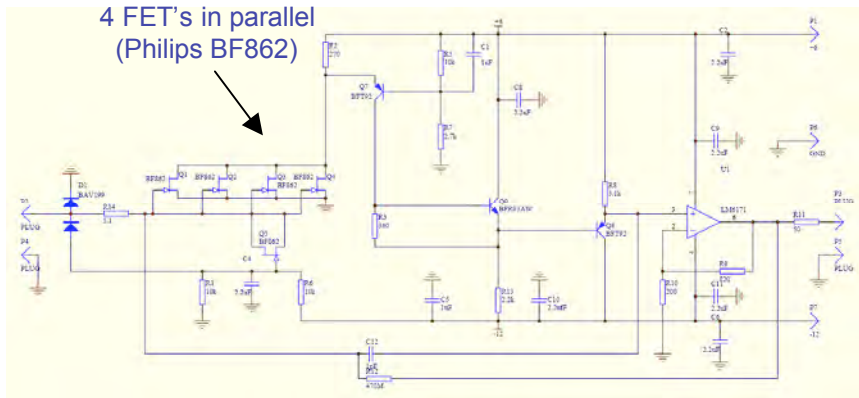
Cables are going through a slot in a UHV flange. The slot is sealed with epoxy resin.

Low noise charge preamp inspired from
C. Boiano et al. IEEE Trans. Nucl. Sci.
52(2004)1931

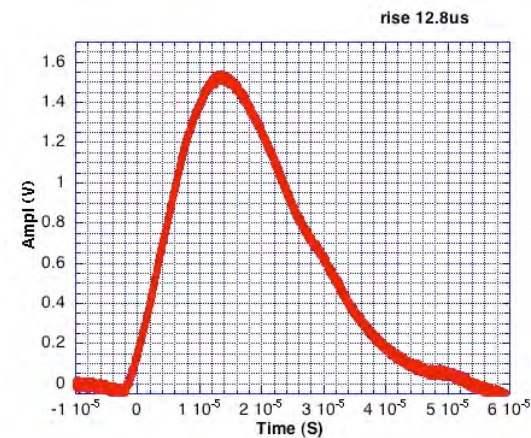
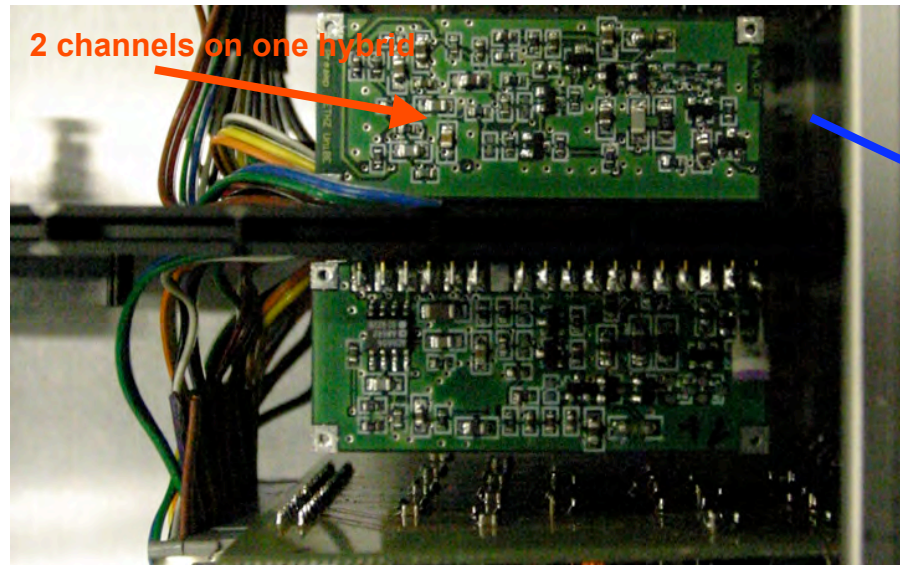
Readout Electronics

Developing A/D conversion and DAQ system:
MHz serial ADC + FPGA + dual memory buffer +
ARM microprocessor

Industrial version being developed
with CAEN

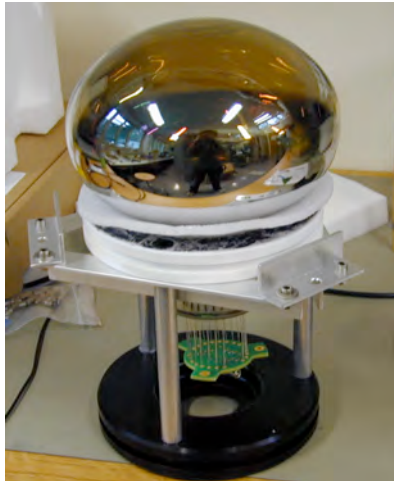


Custom-made front-end charge preamp + shaper
 $G \sim 15\text{mV/fC}$



Signal from double-phase setup

Light readout

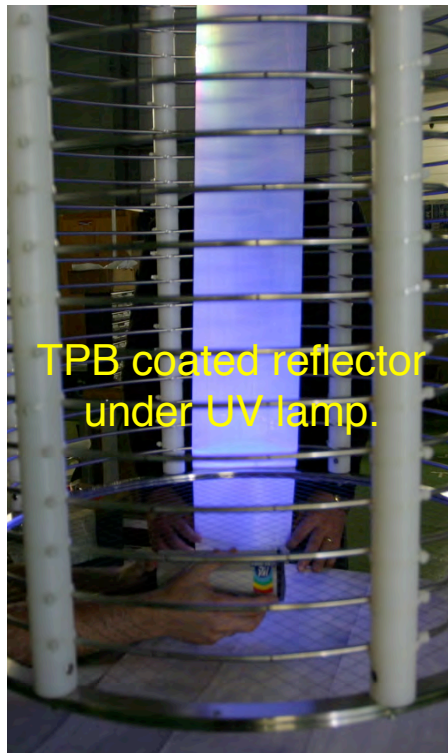


14 low background photomultiplier tubes cover the bottom of the detector

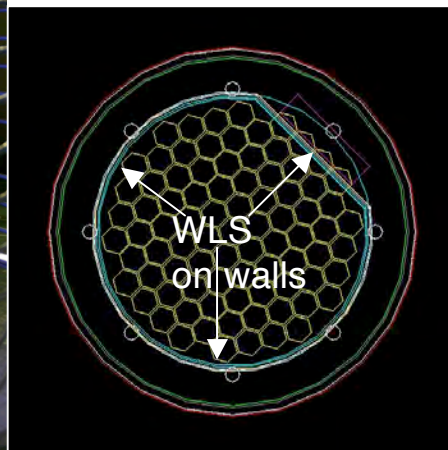
Photomultiplier tube:
Hamamatsu R5912-02MOD
20.2 cm diameter

Wavelength shifter (WLS):
Tetra-Phenyl-Butadiene (TPB)
evaporated on reflector

Reflectivity @430nm ~97%
Shifting eff. 128→430nm >97%



TPB coated reflector under UV lamp.

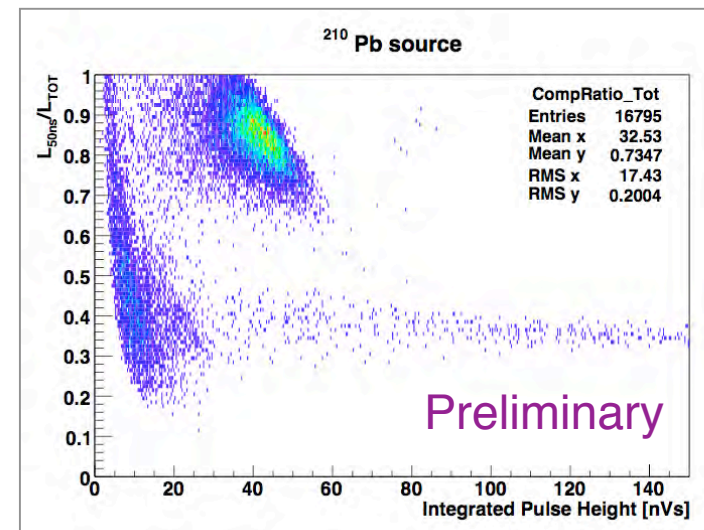


P.Otyugova, ETH Zurich

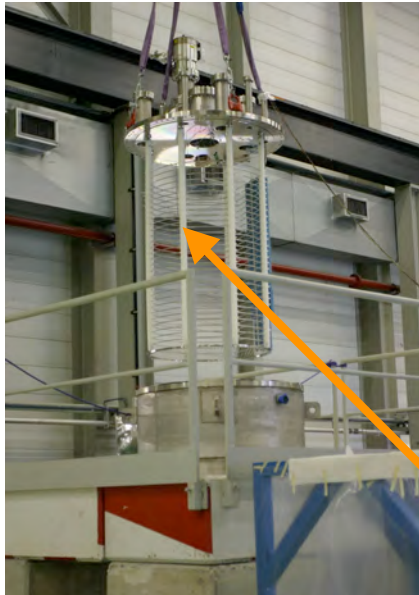


Small test setup

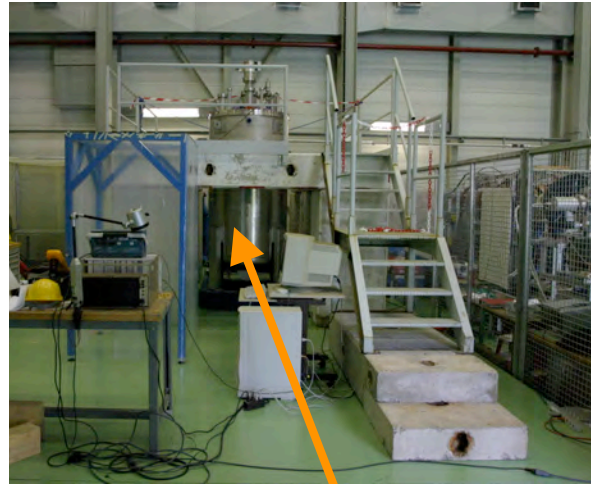
Radioactive source: ^{210}Pb ,
 α 5.3 MeV, β 1.16 MeV
 α and β events are clearly separated.



Detector assembly at CERN



Upper flange



Detector inner part

Side view of the setup



Concrete platform

Greinacher HV system. 210 stages
It has been completed and connected to shaping rings

Cathode mounted on the bottom of the support pillars



Background studies

Background sources:

Neutrons:

From U/Th contaminations of the detector components, muon induced neutrons.

Neutron events look like WIMP-events

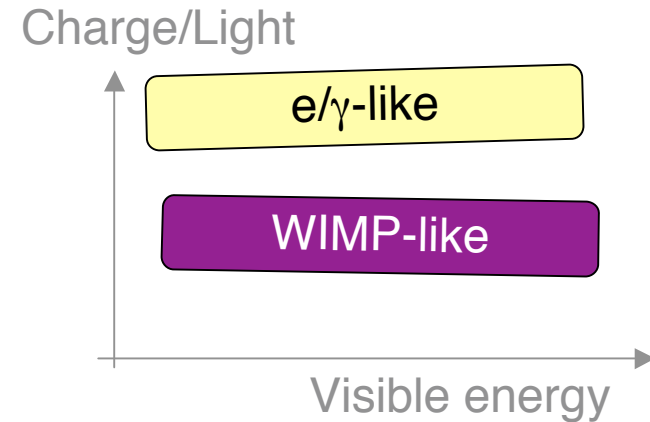
Electrons/Gammas:

From U, Th, K contaminations of detector and surrounding rock.

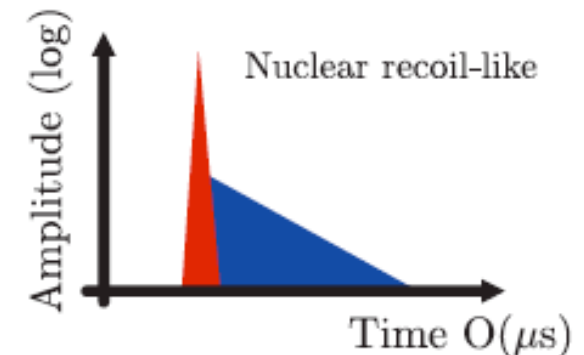
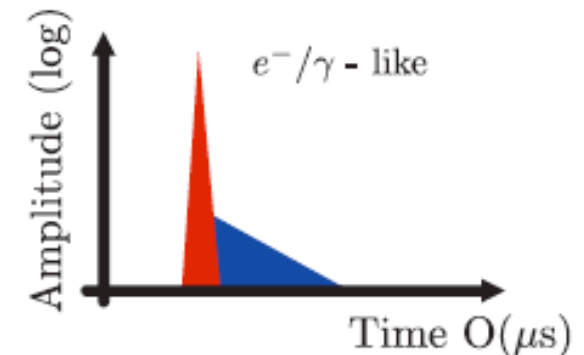
Electron/Gamma events look different from WIMP-events

How can we reject the e/γ background:

- Different light/charge ratios
- Different shape of the scintillation light (ratio fast/slow components).



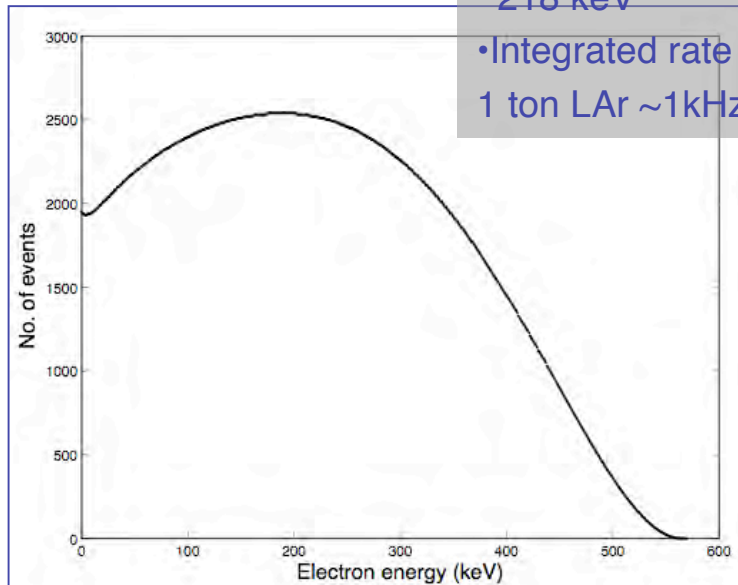
Light:



Ar³⁹ and neutrons backgrounds

Natural argon from liquefaction of air contains small fractions of ³⁹Ar radioactive isotope.

- β -radioactive isotope
- Half life: 269 years
Q=565 keV
- Mean Energy:
218 keV
- Integrated rate in
1 ton LAr ~1kHz



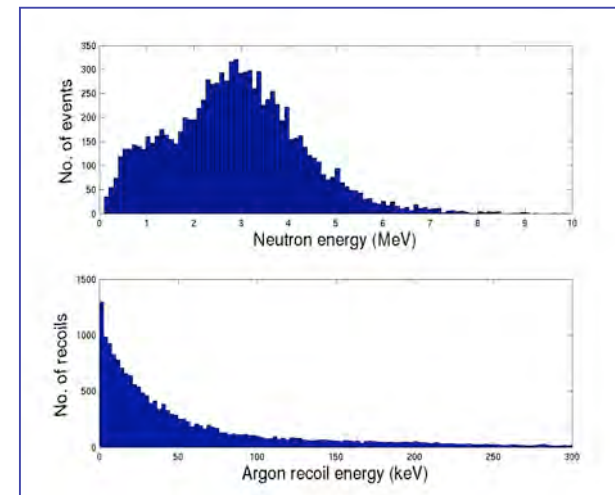
We need:
Rejection power of 10^8
OR
use of ³⁹Ar-depleted argon

Component	n per year	WIMP-like recoils per year
Container	~ 400	~ 30
LEM (std. mat.)	~ 10000	~ 900
LEM (low bg. mat.)	< 20	< 2
14 PMTs (std. mat.)	~ 12000	~ 1000
14 PMTs (low bg. mat.)	~ 600	~ 50

About 55% of the interacting neutrons scatter more than once at the threshold of 30keV.

Less than 10% of the emitted neutrons produce WIMP-like events single recoils, energy $\in [30, 100]$ keV).

The WIMP cross-section is very low, and it will scatter at most once.



ArDM schedule for the near future

- Test of detector in vacuum, at CERN:
High voltage system, purity
Currently in preparation
- Test with gaseous argon, at CERN:
PMTs, high voltage system and small version of LEM plates
Next month
- Test in liquid argon, at CERN:
Recirculation and purification system
Before end of 2007
- Test underground at shallow depth
2008?

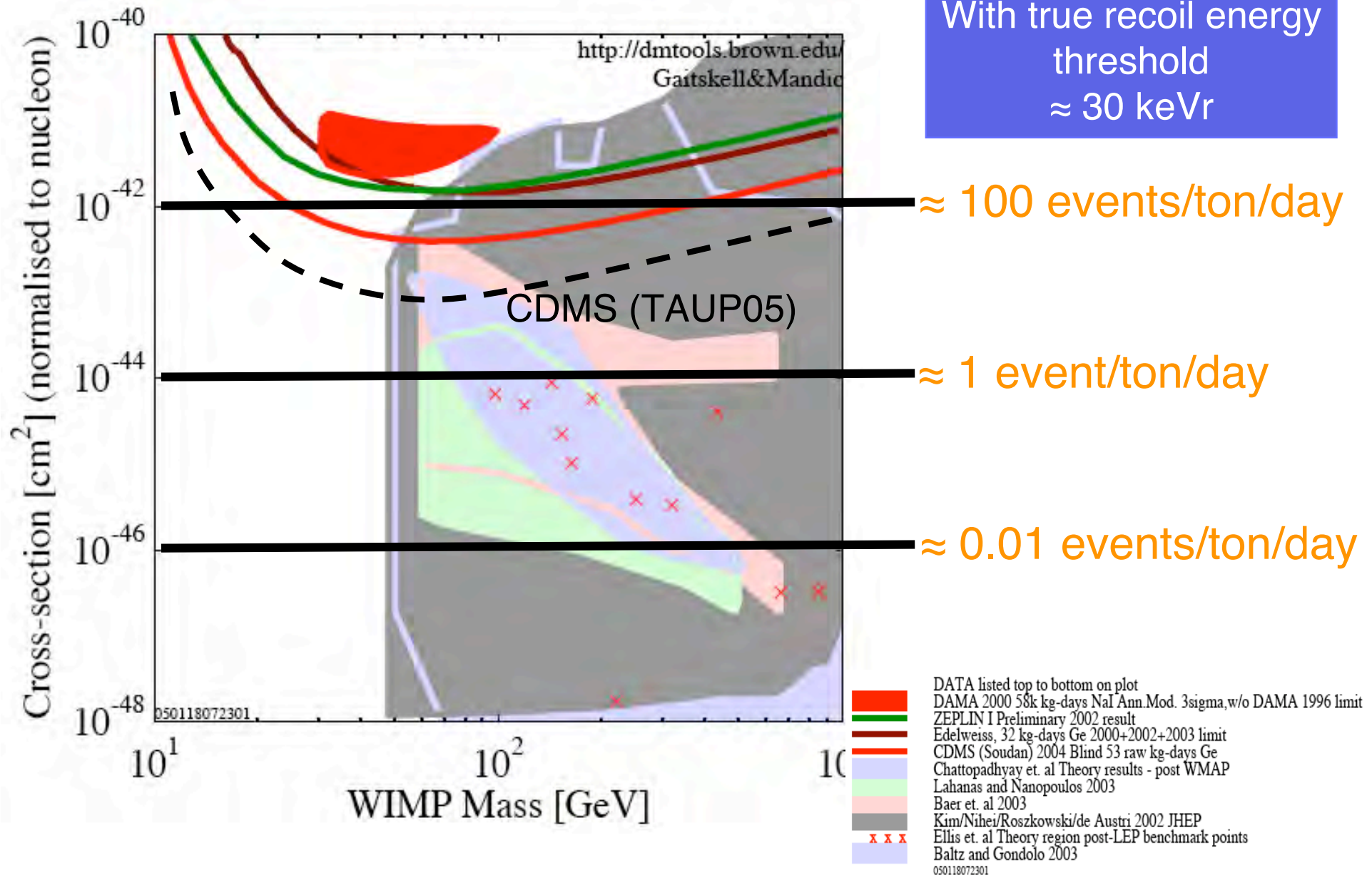
Summary

1. A 1-ton prototype is being assembled at CERN to be run in a first phase above ground (2007).
2. First detector test at CERN are ongoing.
3. Key components:
 - high drift field device
 - LEM-based charge readout
 - argon scintillation light detection system
4. After tests at CERN and at shallow depth will be completed, the detector will be moved to the underground laboratory (presumably to the Canfranc underground laboratory in Spain)
5. The expected sensitivity of the detector will be of the order of 10^{-44} cm^2 (10^{-8} pb), depending on the background rejection power.
6. This technology could be scaled. Detectors of 10 tons and more based on the same technology can be constructed.

Backup slides

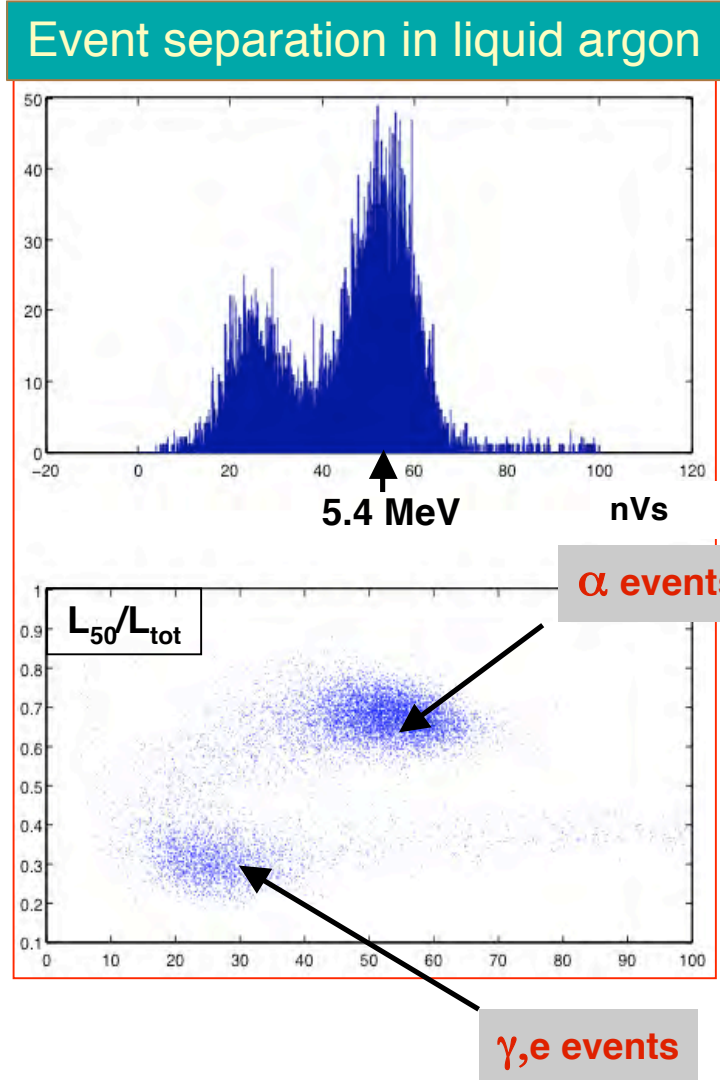
Estimated event rates on argon

With true recoil energy threshold ≈ 30 keVr

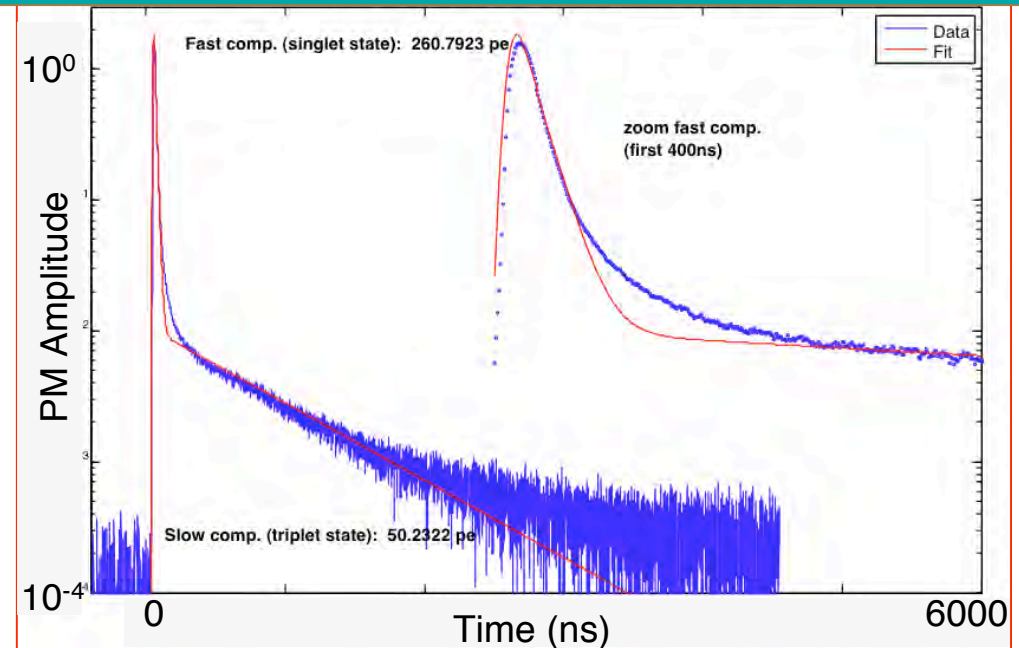


Light measurements in liquid argon

Radioactive source: α (5.4 MeV) + β (Q = 1.163 MeV)

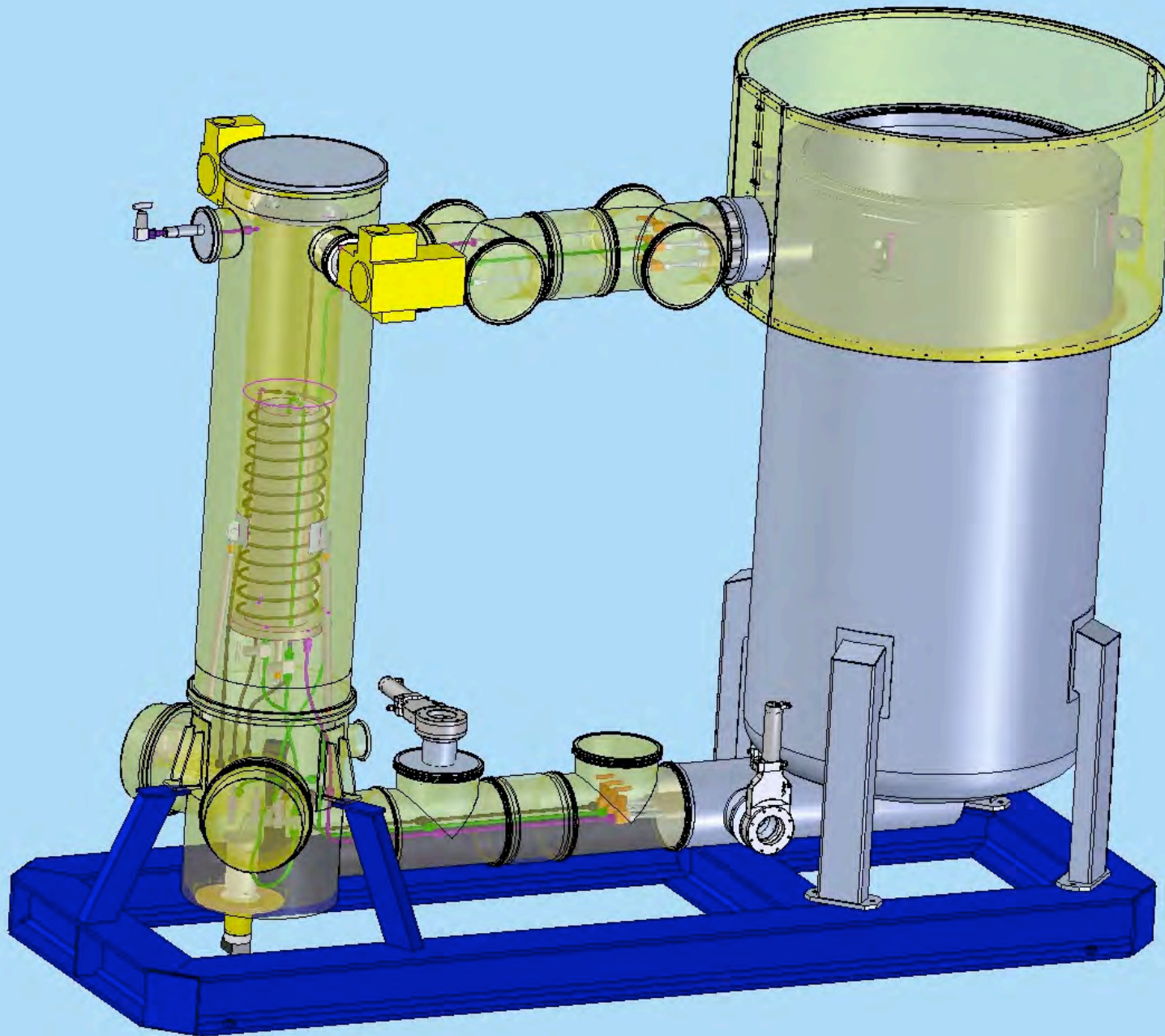


Scintillation light from α in 1200mbar liquid argon



- α events separate well from γ, e events
- Fast and slow light components distinguishable

LAr recirculation system

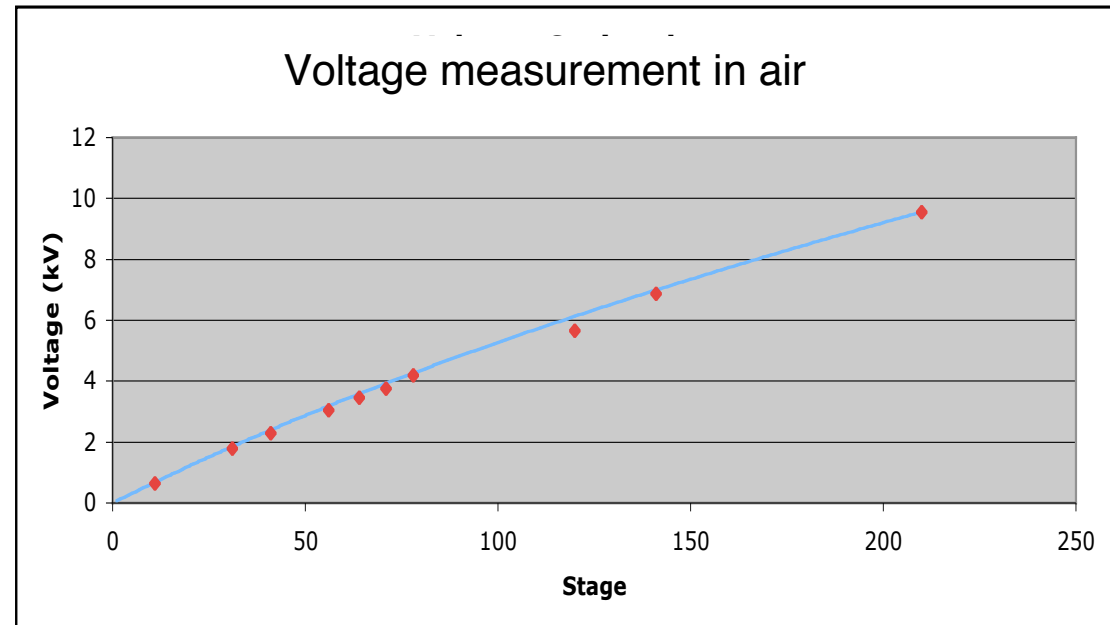


High voltage system

We use a cascade of HV multiplication stages (Greinacher/Cockroft-Walton circuit) directly connected to the field shaping rings

The voltage at the last stage is designed to reach 500 kV, i.e. ≈ 4.17 kV/cm

The Greinacher circuit has been completed and connected to the field shaping rings



Small nonlinearity of the voltage distribution can be corrected with attachments to field shapers

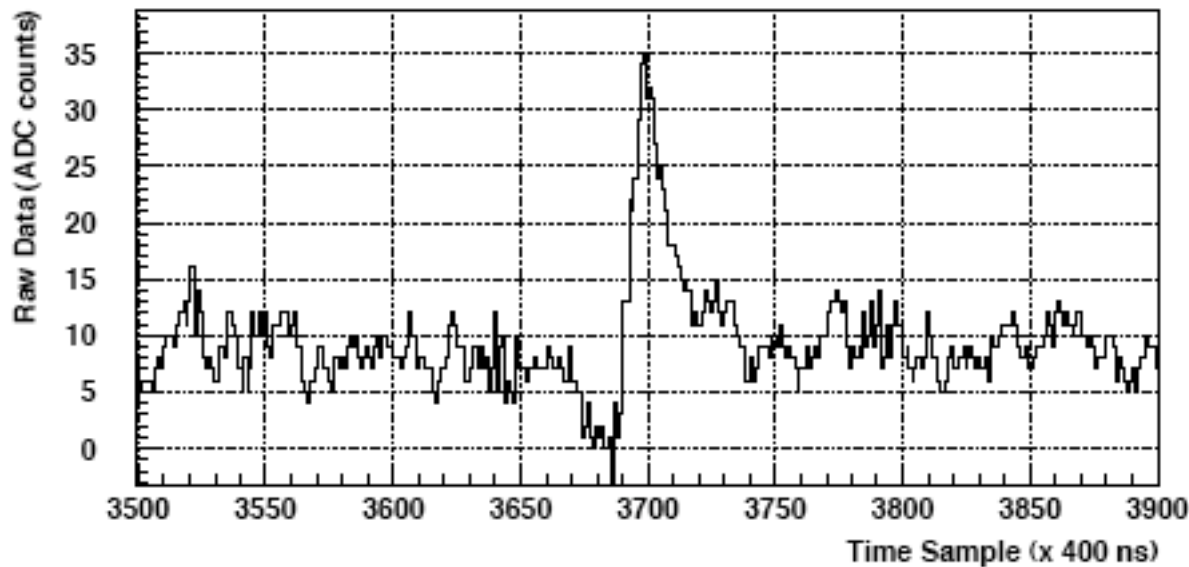
Cathode mounted on the bottom of the support pillars



Liquid Argon TPCs detect the ionization charge to create the image of the event and the scintillation light can be used for triggering or T_0 definition.

To detect the ionization charge in a large noble liquid detector very low noise charge preamplifier is required (challenging, costly).

For example: in ICARUS detector the signal is only 15000 electrons for a minimum ionizing particle track with 3mm wire pitch. In this case to obtain a high signal to noise ratio the equivalent noise charge has to be less than 1000 electrons.



Signal from a MIP recorded In the Induction plane of T600 ICARUS detector.

In 100 kton detector with 20 m drift In a field 1kV/cm the drift time is about 10 ms. With a 2 ms electron lifetime , the 6000 electron/mm signal is attenuated by:

$$e^{-t/\tau} \approx 1/150$$

It is too low for a readout as in ICARUS